

Three-dimensional seismic stability analysis of slopes with linearly increasing undrained shear strength

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(Received January 7, 2025, Revised April 3, 2025, Accepted April 7, 2025)

Abstract. Undrained stability of slopes has long been of interest to the community of geotechnical engineering due to its practical importance. However, the role of the undrained shear strength c_u in the seismic stability and failure geometry of slopes under three-dimensional (3D) condition has not been fully understood. Therefore, this study adopts a modified 3D rotational failure mechanism to develop the upper bound solution to the stability number N_s for slopes with linearly increasing c_u with depth based on the kinematic approach of limit analysis. The seismic force is included using the pseudo-static method. Three types of failure mechanisms, i.e., the toe failure, face failure and below-toe failure are considered to capture the critical condition. Stability charts that cover a wide range of parameters and representative failure surfaces are then developed to illustrate the influence of key factors. Results show that the stability and failure geometry of slopes are significantly influenced by the variation ratio of c_u with depth. The most significant outcome is that an increase in the horizontal seismic coefficient k_h leads to a shallower critical failure surface when the slope with a large variation ratio of c_u is constrained within a narrow width.

Keywords: 3D slope stability; failure mechanism; limit analysis; undrained shear strength

1. Introduction

Seismic hazards can lead to significant consequences, such as landslides and liquefaction of soils (Mase *et al.* 2020, 2021, 2022, 2024), which may induce severe damage to properties, thus assessing seismic slope stability is of concern in geotechnical engineering. Over the past decades, the stability of slopes has been assessed using various numerical methods, such as finite element method (Ukritchon *et al.* 2020, Naeij *et al.* 2021, Satyanaga *et al.* 2022, Petchkaew *et al.* 2023a, 2023b, Wang *et al.* 2024), discrete element method (Bonilla-Sierra *et al.* 2015, Weng *et al.* 2019, Gao *et al.* 2021, Ajith *et al.* 2024), material point method (Bhandari *et al.* 2016, Feng and Xu 2021, Wang *et al.* 2025) and smooth particle hydrodynamics (Ray *et al.* 2019, Zhang *et al.* 2024). Numerical methods have advantages in dealing with slopes with complicated geometric forms and loading conditions, therefore they are extensively adopted for case studies (Bi *et al.* 2019, Azarafza *et al.* 2021, Tran *et al.* 2019, 2023). The limitation of numerical methods lies in the intensive computational effort when carrying out parametric studies. In contrast, analytical approaches including the limit-equilibrium method (Deng *et al.* 2019, Chaudhary *et al.* 2024) and limit

analysis method (Rao *et al.* 2021, Chehade *et al.* 2022) are more convenient for performing parametric studies. The limit equilibrium method is the most frequently used method in engineering practice. Despite its popularity, some limitations still exist in the assumptions with respect to the inter-column forces, the direction of resultant shear resistance and moment equilibrium conditions (Zhou and Qin 2020). The limit analysis method offers a rigorous upper/lower bound for slope stability problems by constructing the kinematically admissible velocity field/statically admissible stress field (Chen 1975). A rigorous upper bound solution to the critical slope height or the factor of safety can be obtained by assuming a reasonable slip surface.

When assessing the undrained stability of slopes or excavations in normally consolidated clays, the undrained shear strength c_u was generally assumed to linearly increase with depth (Hunter and Schuster 1968, Kumar *et al.* 2014). Koppula (1984) assessed the stability of slopes with linearly increasing c_u , where the undrained strength is positive finite on the top surface of the slope. The results demonstrated that the variation ratio of c_u significantly affects the slope stability and the depth of failure surface. Griffiths and Yu (2015) investigated the influence of a firm stratum on the stability and failure surface of slopes with linearly increasing c_u . The zero strength at the crest level was also examined. Some novel results were obtained compared with those of Hunter and Schuster (1968). Li *et al.* (2018) presented a procedure for analyzing the impact of tension crack on the stability of slopes with linearly increasing c_u . It was found that the effect of tension crack on the undrained slope stability is significant under the seismic condition. Griffith and Matin (2020) investigated the failure

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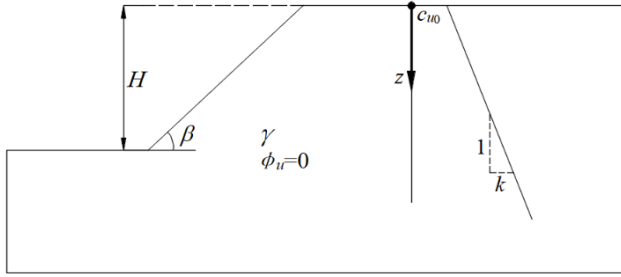


Fig. 2 Problem depiction and notation

4.2 Failure mechanism

As shown in Fig. 3, the failure mechanism becomes a torus for the undrained analysis. Note that when dealing with soils with varied friction angle, the failure mechanism is no longer valid. Three types failure mechanism, i.e., the toe failure, face failure and below-toe failure are considered. The below-toe failure becomes the toe failure when $\beta' = \beta$ (Fig. 3(a)). And the face failure is a special case of toe failure with a smaller slope height H' (Fig. 3(b)). The following equations are formulated based on the below-toe failure mechanism.

The trace of the failure mechanism on the symmetry plane is described by two circular arc AD and AD' , which can be respectively expressed as

$$r = r_0 \tag{3}$$

$$r' = r'_0 \tag{4}$$

The mechanism is generated by rotating a circle defined by the arc given in Eqs. (3)-(4) around the axis passing through the rotation center O . The distance from the rotation center O to the axis of the torus r_m and the radius of the circle R can be expressed as

$$r_m = \frac{r_0 + r'_0}{2} \tag{5}$$

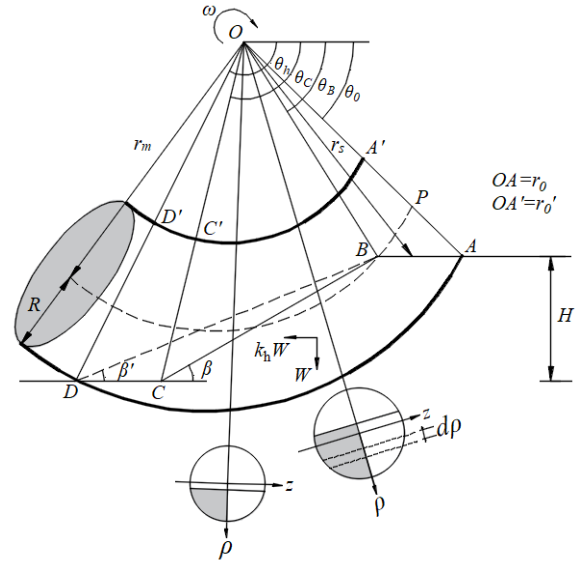
$$R = \frac{r_0 - r'_0}{2} \tag{6}$$

In the polar coordinate system, the contour of the slope can be described as a piece-wise function

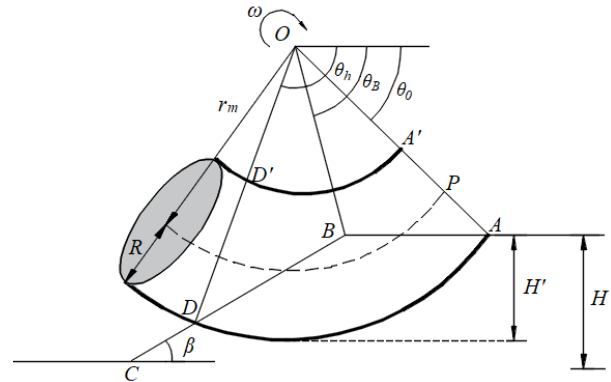
$$r_s(\theta) = \begin{cases} \frac{r_0 \sin \theta_0}{\sin \theta}, & \theta_0 < \theta \leq \theta_B \\ \frac{r_0 \sin(\beta + \theta_h)}{\sin(\beta + \theta)}, & \theta_B < \theta \leq \theta_C \\ \frac{r_0 \sin \theta_h}{\sin \theta}, & \theta_C < \theta \leq \theta_h \end{cases} \tag{7}$$

The z coordinate for the c_u profile can be expressed as

$$z = r_0 (\sin \theta - \sin \theta_0) \tag{8}$$



(a) Below-toe failure



(b) Face failure

Fig. 3 Three-dimensional undrained failure mechanism

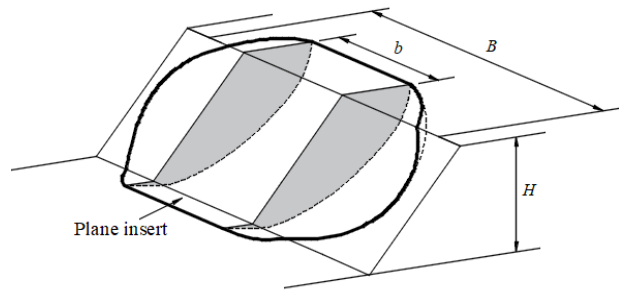


Fig. 4 Mechanism with plane insert

Fig. 4 shows the mechanism for the slope with finite width B modified with a block of width b . When b approaches infinity, the 3D failure mechanism is transformed to the 2D failure mechanism. For a 2D undrained analysis, the depth of the failure mechanism in gentle slopes tends to infinity (Michalowski 2002), hence a firm stratum should be introduced to obtain more rational results. It is noted that this failure mechanism is only valid under undrained condition. If the heterogeneity of the internal friction angle needs to be considered, new failure mechanisms should be developed.

4.3 Solution for the stability number

Based on the failure mechanism depicted in Fig. 3(a), the procedure to determine the stability number N_s will be presented in the following.

For the torus, the work rate done by soil weight can be expressed as

$$W_\gamma^{3D} = 2\omega\gamma \int_{\theta_0}^{\theta_h} \int_{r_s}^r \rho^2 \cos\theta \sqrt{R^2 - (\rho - r_m)^2} d\rho d\theta \quad (9)$$

The vertical acceleration has little effect on slope stability and is ignored herein (Huang *et al.* 2023). The work rate done by the horizontal seismic force can be obtained based on the pseudo-static method

$$W_{k_h}^{3D} = 2\omega\gamma k_h \int_{\theta_0}^{\theta_h} \int_{r_s}^r \rho^2 \sin\theta \sqrt{R^2 - (\rho - r_m)^2} d\rho d\theta \quad (10)$$

where k_h is the horizontal seismic coefficients

The work rate done by soil weight for the plane insert can be expressed as

$$W_\gamma^{2D} = \gamma\omega r_0^3 (f_1 - f_2 - f_3 - f_4) \quad (11)$$

where f_1 - f_4 can be defined as:

$$f_1 = \frac{\sin\theta_h - \sin\theta_0}{3} \quad (12)$$

$$f_2 = \frac{1}{6} \frac{L}{r_0} \sin\theta_0 \left(2\cos\theta_0 - \frac{L}{r_0} \right) \quad (13)$$

$$f_3 = \frac{1}{6} \left[\sin(\theta_h - \theta_0) - \frac{L}{r_0} \sin\theta_h \right] \left[\cos\theta_0 + \cos\theta_h - \frac{L}{r_0} \right] \quad (14)$$

$$f_4 = \frac{1}{2} \left(\frac{H}{r_0} \right)^2 \left[(\cot\beta' - \cot\beta) \cos\theta_0 - \frac{L}{r_0} - \frac{H}{3r_0} (\cot\beta' + \cot\beta) \right] \quad (15)$$

where H/r_0 and L/r_0 can be defined as

$$\frac{H}{r_0} = \sin\theta_h - \sin\theta_0 \quad (16)$$

$$\frac{L}{r_0} = \frac{\sin(\theta_0 + \beta') - \sin(\theta_h + \beta')}{\sin\beta'} \quad (17)$$

The work rate done by the horizontal seismic force for the plane insert can be expressed as

$$W_{k_h}^{2D} = \gamma\omega r_0^3 k_h (f_5 - f_6 - f_7 - f_8) \quad (18)$$

where f_5 - f_8 can be defined as

$$f_5 = \frac{\cos\theta_0 - \cos\theta_h}{3} \quad (19)$$

$$f_6 = \frac{1}{3} \frac{L}{r_0} \sin\theta_0^2 \quad (20)$$

$$f_7 = \frac{1}{6} \left[\sin(\theta_h - \theta_0) - \frac{L}{r_0} \sin\theta_h \right] (\sin\theta_0 + \sin\theta_h) \quad (21)$$

$$f_8 = \frac{1}{2} \left(\frac{H}{r_0} \right)^2 (\cot\beta' - \cot\beta) \left(\frac{2}{3} \frac{H}{r_0} + \sin(\theta_0) \right) \quad (22)$$

Then the total external work rate can be expressed as

$$W_{tot} = W_\gamma^{3D} + W_{k_h}^{3D} + b \cdot (W_\gamma^{2D} + W_{k_h}^{2D}) \quad (23)$$

Due to the rigid body assumption, the internal energy dissipation inside the sliding body is ignored. The internal energy dissipation rate along the failure surface of the torus can be expressed as

$$D^{3D} = 2c_u(x, z) \omega \int_{\theta_0}^{\theta_h} \int_{r_s}^r \rho^2 \frac{R}{\sqrt{R^2 - (\rho - r_m)^2}} d\rho d\theta \quad (24)$$

The internal energy dissipation rate of the plane insert can be expressed as

$$D^{2D} = \int_{\theta_0}^{\theta_h} c_u(x, z) v r_0 d\theta = \omega r_0^2 \int_{\theta_0}^{\theta_h} c_u(x, z) d\theta \quad (25)$$

Then the total internal energy dissipation rate can be expressed as

$$D_{int} = D^{3D} + b \cdot D^{2D} \quad (26)$$

Based on the gravity increase method, the safety factor can be defined as the ratio of the total energy dissipation rate to the total external work rate (Yang and Li 2018)

$$F_s = \frac{D_{int}}{W_{tot}} \quad (27)$$

Therefore, the stability number N_s can be obtained

$$N_s = F_s \cdot \frac{\gamma}{k} \quad (28)$$

To yield the least upper bound solution to N_s , an optimization procedure with respect to the unknown parameters (θ_0 , θ_h , r'/r and β') was coded in MATLAB. The optimization procedure proposed by Chen (1992) was adopted herein. An initial estimate of the unknown parameters can be obtained by the random search method, then the global minimum can be searched by the optimization procedure.

5. Comparison

To validate the robustness of the developed solution and the optimization procedure, the solution in the study is degraded into the homogeneous condition. Herein the calculated dimensionless critical slope heights $\gamma H/c_{u0}$ which can be obtained by equating Eq. (23) to Eq. (26) are

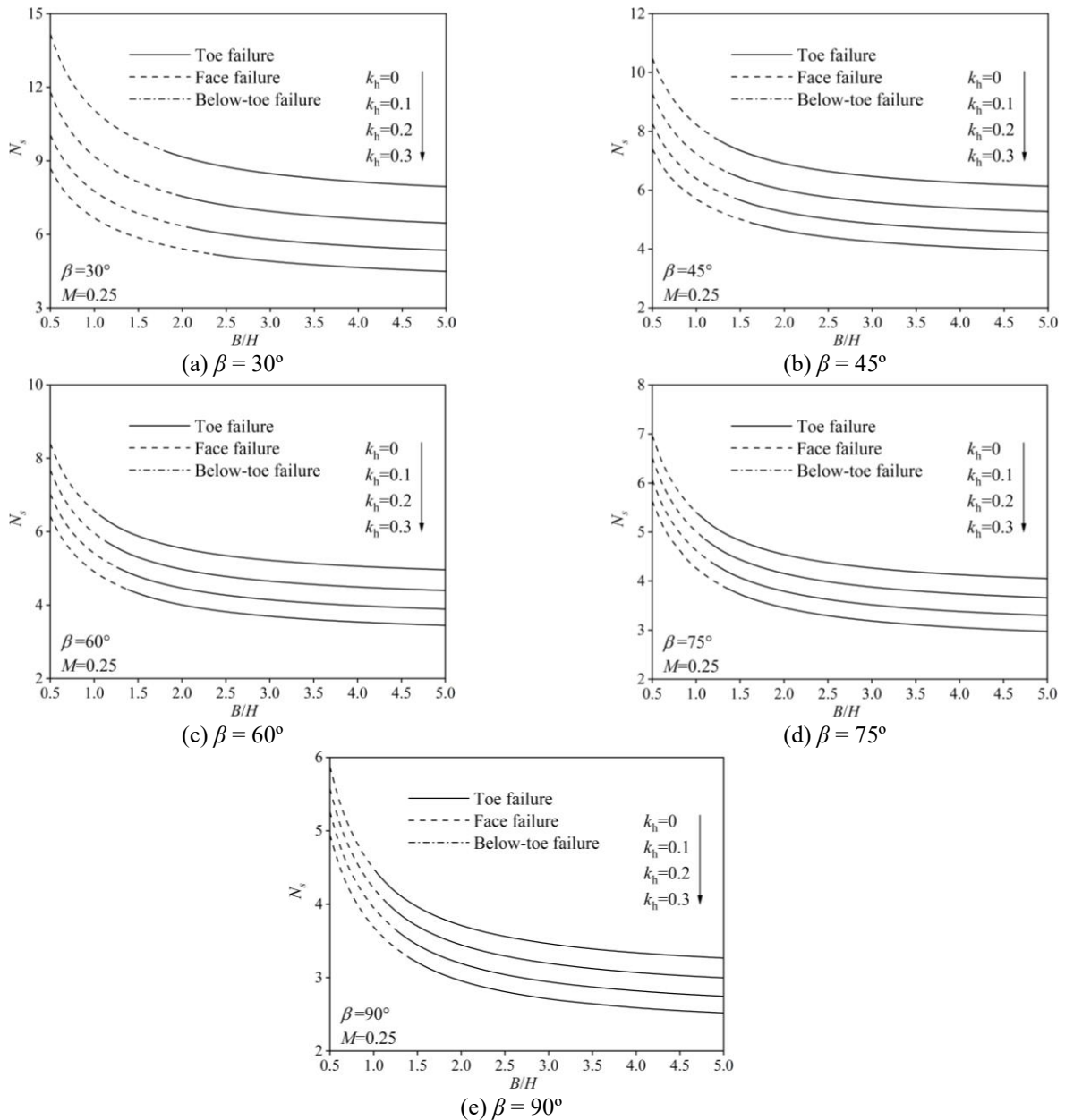


Fig. 5 Variation of stability number N_s with width-to-height ratio B/H for $M = 0.25$

Table 1 Comparison of results

β	B/H	$\gamma H/c_{u0}$	
		This study	Park and Michalowski (2018)
30°	0.5	16.644	16.645
30°	1	10.792 ^a	10.797 ^a
30°	3	7.398 ^a	7.389 ^a
45°	0.5	12.812	12.789
45°	1	9.020 ^a	9.021 ^a
45°	3	6.795 ^a	6.783 ^a
60°	0.5	10.865	10.860
60°	1	7.821	7.821
60°	3	6.022	6.022

^a: Below-toe failure

compared with those obtained by Park and Michalowski (2018), as shown in Table 1. It is clear that the present solution agrees well with the result of Park and Michalowski (2018).

The case of linearly increasing undrained shear strength with $c_u = 0$ ($M = 0$) at the slope crest level under 2D condition was examined by Griffiths and Yu (2015). The present solution will become 2D solution when b approaches infinity, then the stability number N_s obtained using the proposed approach are 8.511, 4.106 and 2.563 for $\beta = 15^\circ, 45^\circ$ and 75° respectively. These results are identical with those obtained by Griffiths and Yu (2015) as indicated in Fig. 7 in their original work. Then the validity of the solution and the optimization procedure is confirmed.

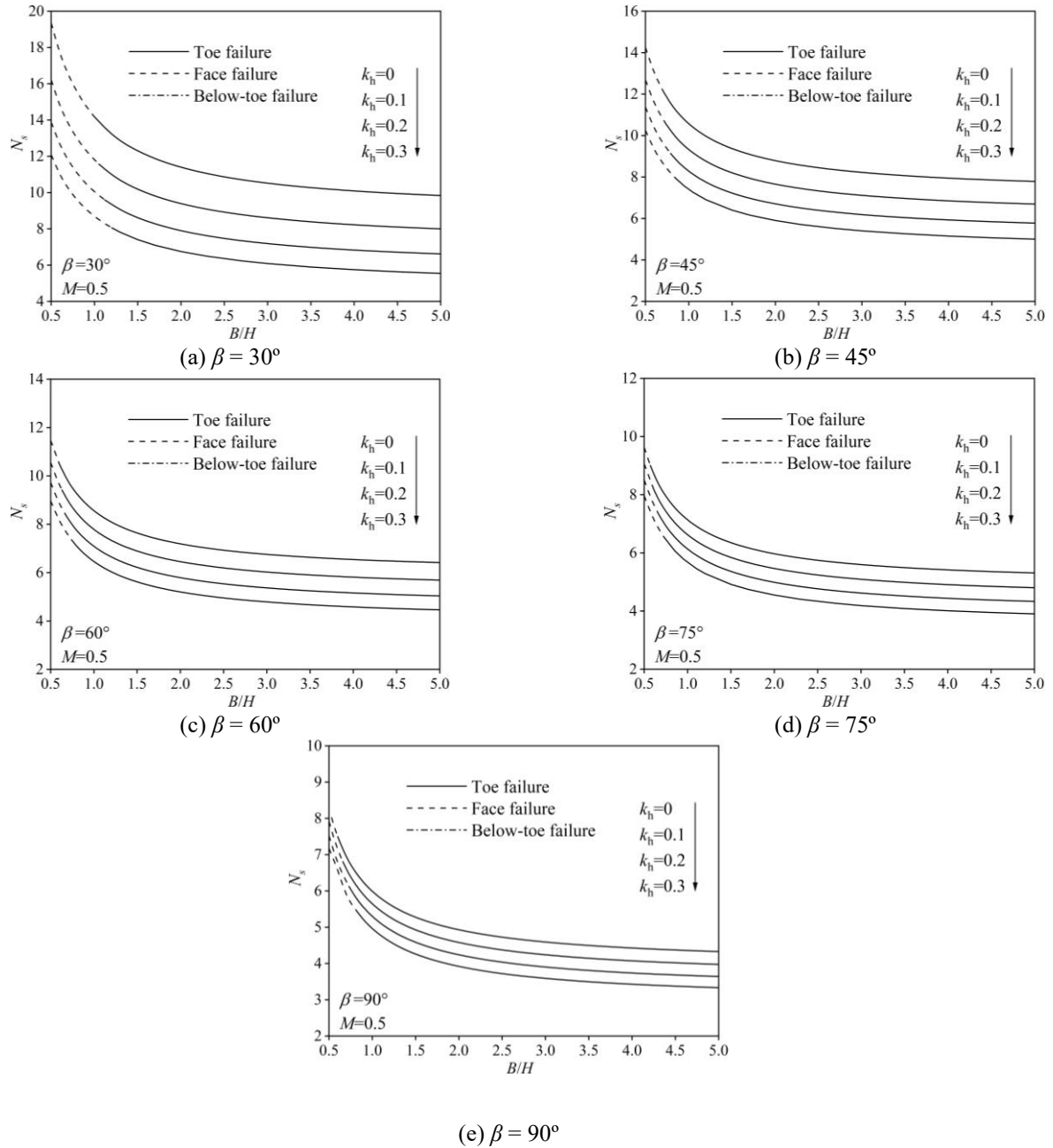


Fig. 6 Variation of stability number N_s with width-to-height ratio B/H for $M = 0.5$

6. Results and discussions

To capture the influence of key parameters, i.e., the c_u profile parameter M , horizontal seismic coefficient k_h and width-to-height ratio B/H on the slope stability and failure geometry, a series of charts will be presented in the following section, both the stability number and the failure mechanism are indicated in the figures (solid lines for the toe failure mechanism, dashed lines for the face failure mechanism and dash-dotted lines for the below-toe failure mechanism).

6.1 Slope stability analysis

The obtained stability charts are presented in Figs. 5-8.

The plots cover a wide range of slope angel β (30° , 45° , 60° , 75° and 90°), the width-to height ratio B/H varies from 0.5 to 5.

From Figs. 5-8, it can be concluded that as B/H increase from 0.5 to 5, the stability number N_s decreases sharply at early stage, and then the decreasing rate is reduced. Moreover, the increase of M gives rise to an increasing N_s , slopes with smaller B/H ratio are more sensitive to the variation of M . For instance, for a 30° slope with $k_h = 0$, by increasing M from 0.25 to 0.5, the stability number N_s increases by 36.6% for $B/H = 0.5$, while an increase of 23.8% is observed for $B/H = 5$. Manifestly, an increase in k_h leads to a decrease in N_s , and gentler slopes are more sensitive to the variation of k_h . As shown in Figs. 5(a) and 5(d), an increase in k_h from 0 to 0.3 leads to a reduction of

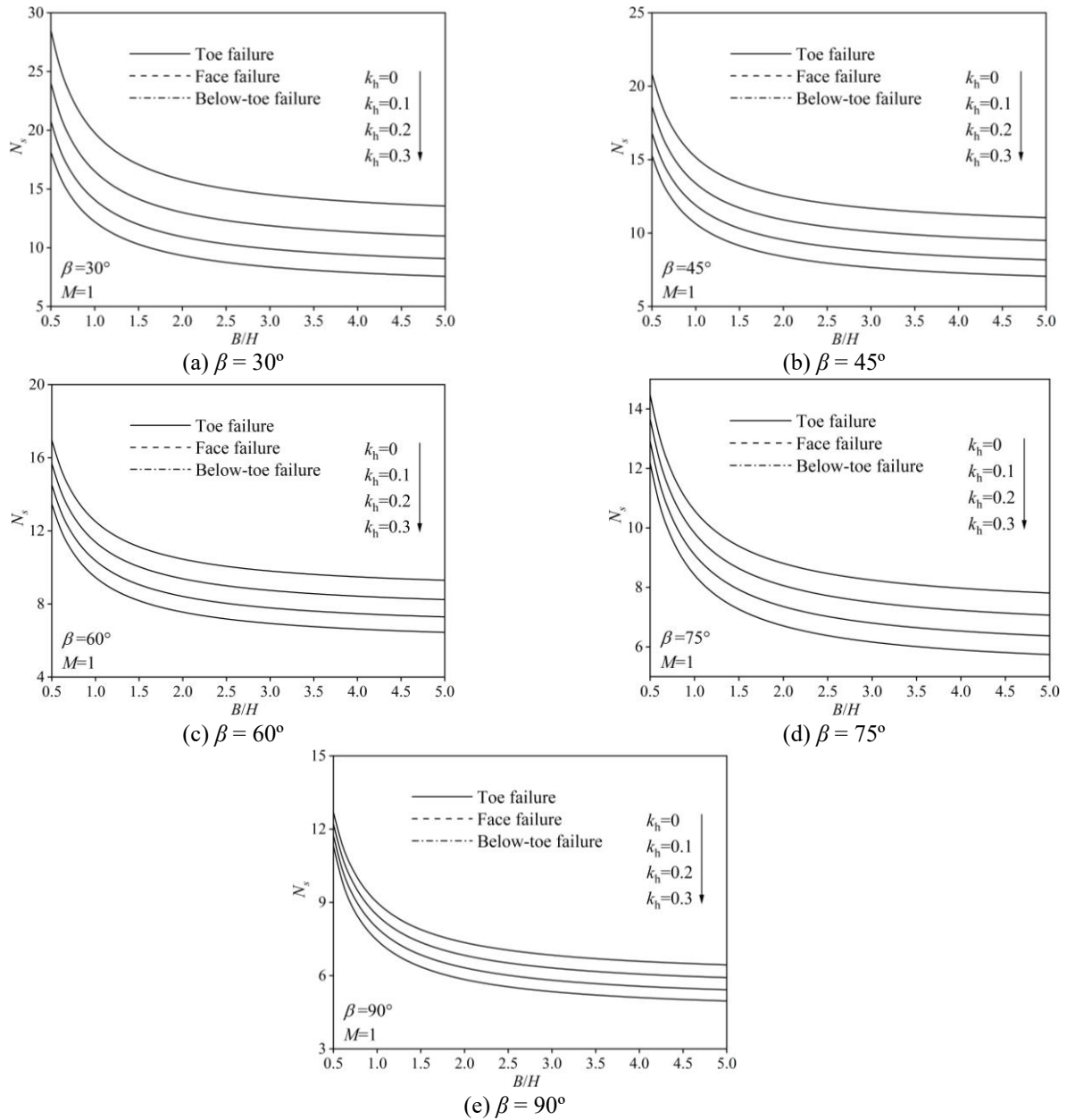


Fig. 7 Variation of stability number N_s with width-to-height ratio B/H for $M = 1$

38% and 15% in N_s for 30° and 90° slopes respectively with $B/H = 0.5$.

As shown in Figs. 5 and 6, when $M = 0.25$ and 0.5 , with the increase of B/H , the failure mode changes from the face failure to the toe failure. An interesting phenomenon can be observed that an increase in k_h leads to more B/H values involved in the face failure mechanism. This phenomenon will be explained later in conjunction with the shape of the failure surface.

Additionally, from the charts, it can be found that the below-toe failure only occurs in gentle slopes when M is large as shown in Fig. 8(a), while the face failure occurs for a wide range of slope angles when M is small. Herein the largest value for M is 2, it is expected that with the further increase of M , the soil strength becomes more uniform, then the failure trends are similar to those observed in

homogeneous slope (Gao *et al.* 2013), and thus for a large M value, the below-toe failure is produced in gentle slopes.

6.2 Failure geometry

Selected failure geometries in the symmetry plane for 60° slopes are presented in Fig. 9, the effects of M , B/H and k_h on the critical failure geometries are illustrated in the figure. As can be seen from Fig. 9(a), with the increase of k_h , the failure geometry becomes shallower, while the extent of the failure surface in the horizontal direction becomes wider. This phenomenon is not common since the seismic force tends to make the critical failure surface deeper from the traditional perspective. But it makes sense because when M is small, the variation ratio of c_u in the depth direction is large, then the strength of the upper layer is

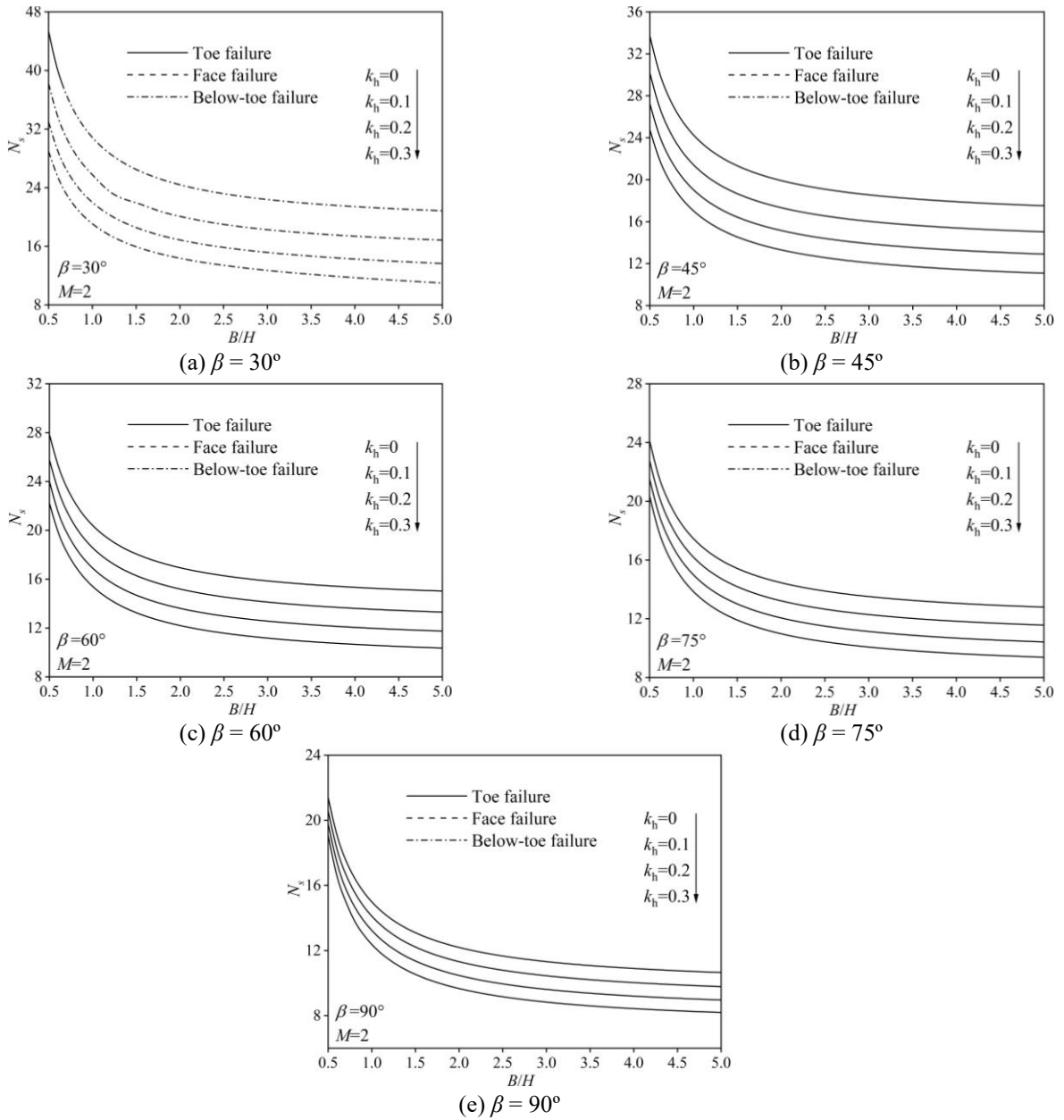


Fig. 8 Variation of stability number N_s with width-to-height ratio B/H for $M = 2$

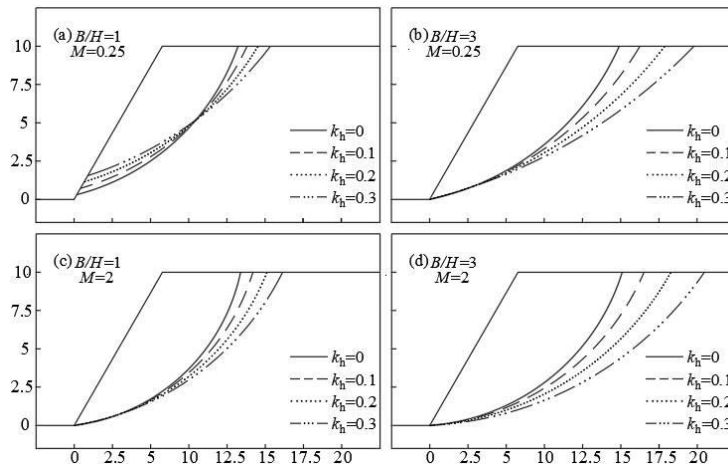


Fig. 9 Failure geometries in the symmetry plane

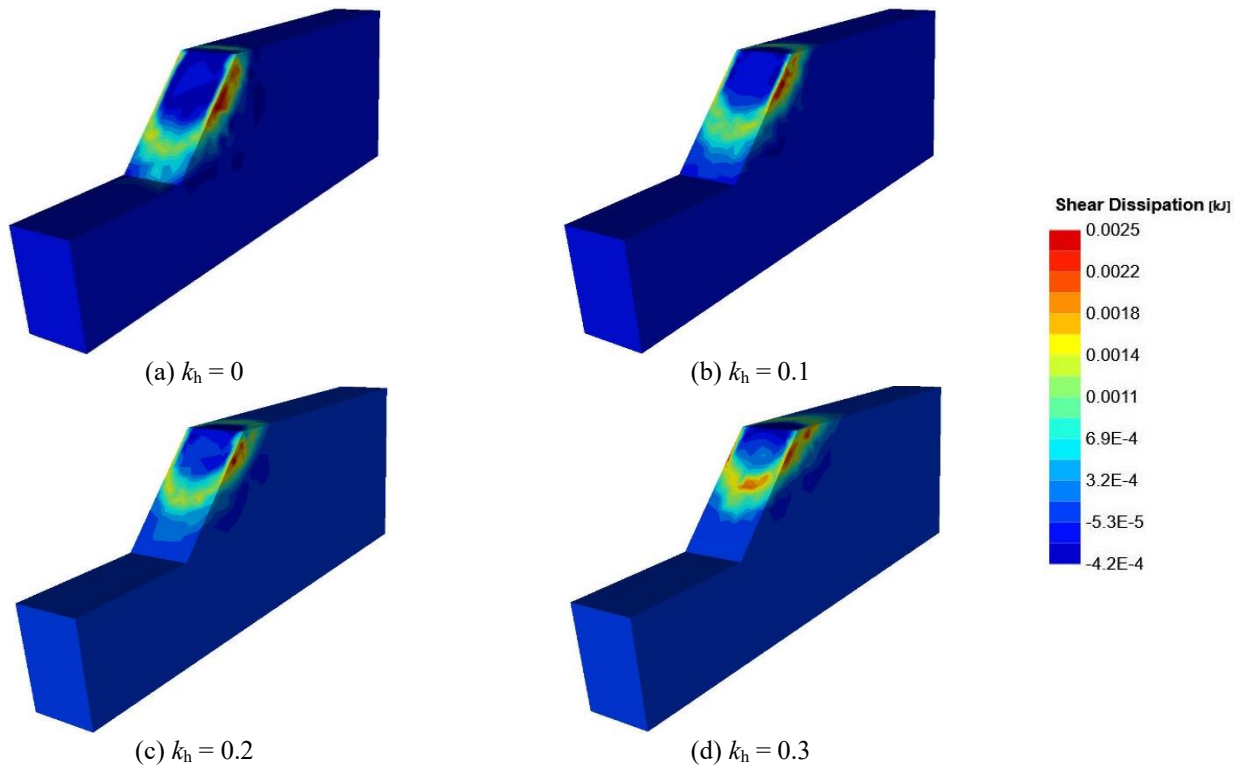


Fig. 10 Three-dimensional failure mechanisms for a slope with $\beta = 45^\circ$, $B/H = 2/3$ and $M = 0.25$

much weaker than that of the lower layer. In this scenario, compared with a deeper failure surface, a shallower failure surface is much more easily to reach the limit state under the seismic condition. Numerical simulations were also carried out based on the FELA to further verify this conclusion. The detailed numerical scheme can be found in Li *et al.* (2009) and Hong-in *et al.* (2023). Fig. 10 shows the three-dimensional failure mechanisms for a 45° slope with $B/H = 2/3$ and $M = 0.25$ under different horizontal seismic coefficient k_h . It can be seen that with the increase of k_h , the failure mechanism becomes shallower. It is noted that the obtained failure surface gives an estimate of the incipient motion of the slope failure, if a stronger seismic excitation occurs after the initial failure, the secondary failure should be considered (Qin *et al.* 2018).

The aforementioned conclusion holds only when both B/H and M are small. For most cases considered, as illustrated in Figs. 9(b)-9(d), an increasing k_h makes the failure geometry deeper and wider. Additionally, comparing Figs. 9(a)-9(d), it is seen that with the increase of M , a larger failure block can be observed under both static and seismic condition. This can be ascribed to the fact that an increasing M indicates a decreasing variation ratio of c_u in the depth direction, namely the strength of soils away from the slope top surface are reduced, which makes more soils prone to failure.

As the stability number and failure mechanism are significantly affected by M and B/H , in engineering practice, using the traditional toe failure mechanism may lead to an erroneous estimate of seismic slope stability, which in turn leads to an inappropriate reinforcement scheme and reinforcement location. It clearly shows the

importance of using extended failure mechanisms rather than the toe failure mechanism only in the seismic design for slopes.

7. Conclusions

This study conducts an analysis on the 3D seismic stability of slopes with linearly increasing undrained shear strength c_u . Based on the kinematic approach of limit analysis and the pseudo-static approach, the upper bound solution to the stability number N_s is obtained. The effects of c_u profile parameter M , horizontal seismic coefficient k_h and width-to-height ratio B/H on the stability and failure geometry of slopes are investigated. The key findings are outlined as follows:

- The increase of M gives rise to an increasing N_s . Slopes with smaller B/H are more sensitive to the variation of M . An increase in k_h leads to a decrease in N_s , and gentler slopes are more sensitive to the variation of k_h .
- The failure geometry is significantly affected by M and k_h . The below-toe failure can only be observed in gentle slopes when M is large, while the face failure is produced for a wide range of slope angles when M is small.
- The width-to-height ratio B/H plays a key role for soils with a large gradient in c_u in the depth direction, i.e., small value of M . A novel finding is drawn that an increase in k_h makes the failure surface shallower when B/H is small.
- This research adopted an analytical method to evaluate the 3D undrained slope stability. Further work should focus on the development of efficient methods for more complex 3D problems, to model the unsaturated conditions, complex

loading conditions and the heterogeneity of internal friction angle.

Acknowledgments

The research described in this paper was financially supported by the National Natural Science Foundation of China (Project No.: 51208301). The authors wish to express their gratitude for the above financial support.

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