

# Nonlinear vibration analysis of FG porous shear deformable cylindrical shells covered by CNTs-reinforced nanocomposite layers considering neutral surface exact position

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**Abstract.** This paper presents nonlinear vibration analysis of a composite cylindrical shell. The core of the shell is made of functionally graded (FG) porous materials and layers is fabricated of carbon nanotubes (CNTs) reinforced nanocomposites. To increase the accuracy of results, neutral surface position is considered. First-order shear deformation theory is used as displacement field to derive the basic relations of equation motions. In addition, von-Karman nonlinear strains are employed to account geometric nonlinearity and to enhance the results' precision, the exact position of the neutral surface is considered. To governing the partial equations of motion, the Hamilton's principle is used. To reduce the equation motions into a nonlinear motion equation, the Galerkin's approach is employed. After that the nonlinear motion equation is solved by multiple scales method. Effect of various parameters such as volume fraction and distribution of CNTs along the thickness directions, different patterns and efficiency coefficients of porous materials, geometric characteristics and initial conditions on nonlinear to linear ratio of frequency is investigated.

**Keywords:** carbon nanotube reinforcements; cylindrical shells; nonlinear vibration; porous materials; von-Karman strains

## 1. Introduction

Different shapes of shells, including cylindrical, spherical, conical, and ellipsoidal, have distinct applications and advantages based on their geometrical properties (Gao *et al.* 2022, Wang *et al.* 2023). Cylindrical shells are a type of structural component that finds wide-ranging applications in various fields due to their specific geometrical properties and advantages. Let's explore different shapes of shells, with a focus on cylindrical shells, and discuss their applications and advantages. Cylindrical shells are essentially hollow, tube-like structures with a circular or cylindrical cross-section. They consist of two main components: the outer curved surface and the inner curved surface. These surfaces are typically formed by rolling or bending flat sheets of material. They are commonly used in the construction of pressure vessels, such as boilers and tanks, where they are designed to withstand internal pressure (Han *et al.* 2018, Tian *et al.* 2024, Wang *et al.* 2017). Their curved shape allows them to distribute pressure evenly, making them suitable for containing gases or liquids at high pressures. They are also used in the construction of pipelines for transporting various fluids,

including oil, gas, and water. Their streamlined shape reduces fluid resistance, making them energy-efficient for long-distance transport. Cylindrical shells are used as structural components in buildings, bridges, and other engineering structures. They offer high strength-to-weight ratios, allowing for efficient load-bearing capacity. In aerospace engineering, cylindrical shells are used in the construction of aircraft fuselages and rocket casings. Their aerodynamic shape reduces drag and enhances flight performance. Cylindrical tanks are widely used for the storage of liquids and bulk materials. Their cylindrical shape maximizes storage capacity and minimizes material wastage. Consequently, understanding the characteristics of these shell shapes is crucial in engineering and design to select the most appropriate shape for a given application (Dehghan *et al.* 2020, Kalamkarov *et al.* 2006, Kiarasi *et al.* 2021, Mirjavadi *et al.* 2020a, b, Ni *et al.* 2023, Wang *et al.* 2019).

Porous materials, by definition, contain voids or empty spaces within their structure. These voids can range from nanometers to millimeters in size and may be interconnected or isolated. These materials are prized for their unique properties and have a diverse range of applications across various fields. Porosity is the defining characteristic of porous materials. It quantifies the ratio of void volume to total volume and is usually expressed as a percentage. High porosity indicates a larger fraction of empty space within the material. The pore size distribution is a crucial specification, describing the range of pore sizes within the material. This distribution influences properties

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like permeability, surface area, and the material's ability to adsorb or release substances. Permeability measures the ease with which fluids or gases flow through the porous material. It depends on both porosity and pore connectivity. Materials with high permeability are valuable for filtration, drainage, or fluid transport applications. Porous materials often exhibit a high surface area due to their numerous pores and internal surfaces. This high surface area is advantageous for applications involving adsorption, catalysis, or gas storage. Bulk density represents the mass of a porous material per unit volume, considering both the solid material and the void spaces. It's a critical specification for structural materials. The mechanical strength of porous materials varies based on composition and fabrication methods. Some are designed to be lightweight with low mechanical strength, while others are engineered for high strength. Porous materials can have unique thermal properties, such as low thermal conductivity, making them suitable for use as insulators or thermal barriers. Chemical reactivity can be tailored by modifying the surface properties of porous materials. Functionalization enhances their ability to adsorb or react with specific substances. In terms of electrical properties, some porous materials exhibit electrical conductivity due to conductive particles within the pores, while others may be insulating or semiconducting. Porous materials have diverse applications in various fields, including catalysis, gas adsorption, water purification, air filtration, biomedical uses (like tissue engineering and drug delivery), energy storage (e.g., batteries and supercapacitors), and construction (for insulation and acoustical control). Researchers and engineers continually explore and develop new porous materials, tailoring their properties for emerging applications in science and technology. These materials, with their versatile characteristics, contribute significantly to advancements across multiple domains (Al-Osta *et al.* 2021, Arefi *et al.* 2019a, Arshid *et al.* 2023a, b, Berghouti *et al.* 2019, Cong *et al.* 2018, Mohammad-Rezaei Bidgoli and Arefi 2019, Mousavi *et al.* 2021, Quang *et al.* 2021, Soleimani-Javid *et al.* 2021, Tahir *et al.* 2021, Zeng *et al.* 2019).

Nanocomposite materials are a class of advanced materials that combine a matrix material with nanoscale reinforcements. These reinforcements can be in the form of nanoparticles, nanofibers, or nanosheets, and they are typically dispersed within the matrix at the nanometer scale (Su *et al.* 2023, Sun *et al.* 2024). The integration of nanoscale reinforcements imparts unique properties to these materials. Nanocomposites feature nanoparticles with dimensions typically ranging from 1 to 100 nanometers. The uniform dispersion of these nanoparticles within the matrix is critical for enhancing material properties. Nanocomposites have a matrix material that can be a polymer, metal, ceramic, or a combination of these. The choice of matrix material depends on the desired properties and application. The nanoscale reinforcements can be of various types, including carbon nanotubes (CNTs), graphene, clay nanoparticles, metal nanoparticles, and more. Each type imparts specific properties to the nanocomposite. Many nanocomposites are lightweight and

have a high strength-to-weight ratio, making them ideal for aerospace and automotive applications. Nanocomposite materials find applications in a wide range of industries, including aerospace, automotive, electronics, construction, biomedical, and energy storage. They are used in structural components, coatings, sensors, drug delivery systems, and more (Arefi *et al.* 2018a, Bi *et al.* 2021, Djilali *et al.* 2022, Forsat *et al.* 2021, Paul *et al.* 2013, Safaei *et al.* 2019, Shen *et al.* 2017b, Sobhani *et al.* 2022). The key outstanding characteristics of the FG porous core subjected to compression between CNT-reinforced composite layers, compared to traditional functionally graded materials (FGM), are explained in the following. The addition of CNT reinforcement to the composite layers surrounding the FG porous core provides significant improvements in tensile strength, compressive strength, and impact resistance compared to FGMs without CNT reinforcement (Ninh *et al.* 2023, Ninh and Bich 2016, Tien *et al.* 2021). The CNTs help to effectively transfer and distribute loads through the composite layers (Taj *et al.* 2022, 2023a, b). The FG porous core allows for a more continuous and controllable property gradient to be achieved, compared to the step-wise property changes typical of laminated FGM structures (Babaei *et al.* 2019, 2022, Babaei and Asemi 2022). This enables better optimization of the stress/strain distribution and overall mechanical performance. The porous nature of the core material allows for enhanced energy dissipation and absorption capabilities during compression, impact, or other dynamic loading events. This can make the material more effective for applications like lightweight armor or impact protection. The ability to independently vary the composition and properties of the porous core and the CNT-reinforced composite layers provides greater design flexibility compared to monolithic FGMs. This allows the material properties to be tuned for specific functional requirements. Depending on the specific materials used, the composite structure may also exhibit improved thermal insulation or heat dissipation properties compared to homogeneous FGM counterparts (Arefi *et al.* 2019b, Babaei *et al.* 2024, Bidgoli *et al.* 2022, Khoddami Maraghi and Arshid 2024, Mohammad-Rezaei Bidgoli and Arefi 2023b).

Numerous researches have been conducted by scholars to find out mechanical performance of engineering structures, especially shells that are made from porous and nanocomposite materials (Babaei *et al.* 2020, Bi *et al.* 2023, Kiarasi *et al.* 2020, Shen *et al.* 2024). For instance, in earlier studies, Das (1964) and Dong (1968) investigated the vibrations of orthotropic and laminated orthotropic cylindrical shells. Sivadas and Ganesan (1991) focused on the vibrations of orthotropic shells with variable thickness. Vronay and Smith and (1970) explored the vibrational behavior of finite-length circular cylindrical shells, while Qatu and colleagues (2004, 2010) conducted research on laminated composite shells and plates, examining their dynamic response. Liu *et al.* (2012) discussed the vibrations of orthotropic cylindrical shells. Additionally, various studies have been conducted in the literature on reinforced composites. Shen (2013) delved into FGCNT-reinforced composite plates, considering thermal buckling and post-buckling effects. They factored in temperature variations

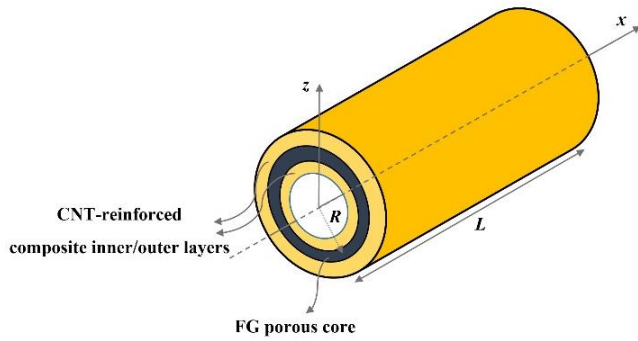


Fig. 1 Schematic of the under-evaluation sandwich composite shell

affecting single-walled carbon nanotubes (SWCNTs) properties and employed a multi-scale approach to obtain their results. Wang and Shen (2012) investigated the impact of thermal conditions on the dynamic response of CNT-reinforced composite plates, considering two CNTs distribution patterns and observing that increased CNTs volume fractions led to higher natural frequencies. Malekzadeh and Shojaee (2013) explored the buckling of quadrilateral laminated plates with CNT-reinforced composite layers. Aragh *et al.* (2012) employed the 2D generalized differential quadrature method to analyze the vibration of continuously graded CNT-reinforced cylindrical panels under various boundary conditions. Jam *et al.* (2012) utilized 3D elasticity to study the vibration of CNTRCs panels, taking into account the aspect ratio and waviness of CNTs. Mohammadimehr and co-authors (2016, 2019) investigated CNT-reinforced composite extensively, considering them as face sheets in sandwich structures. They examined temperature-dependent properties, small-scale behavior, and the impact of electric and magnetic fields on CNT-reinforced composite, presenting their findings and discussing the effects of different parameters. Yas *et al.* (2013) conducted a vibrational analysis of CNT-reinforced composite panels using 3D elasticity, comparing different CNTs distribution patterns within the panels. Shen and his co-authors in different works (Shen and Xiang 2014, Shen 2011, Shen *et al.* 2017a, Shen and Xiang 2013) contributed research on post-buckling and nonlinear vibration of CNT-reinforced composite shells and panels, considering both thermal and mechanical buckling. Shen and Xiang (2014) studied the vibrations of CNT-reinforced composite cylindrical panels with large amplitude deflections, employing the higher-order shear deformation theory for accurate results. This theory is particularly useful for thick structures due to its consideration of shear deformation effects, albeit resulting in more complex equations. Finally, Arshid *et al.* (2022) recently conducted an analysis of higher-order hygro-magneto-electro-thermo-mechanical behavior in FG nanocomposite cylindrical shells with piezo-electro-magnetic layers, with a focus on linearity in their work.

Informed by previous research efforts, this paper aims to advance the understanding of sandwich cylindrical shells and explores their nonlinear vibrational characteristics. It's worth noting that within the realm of published studies, this research constitutes a pioneering analysis of such structures.

To align our findings more closely with practicality and harmonize them with experimental studies, the FSDT is employed to model various arbitrary points within the structure. The core of this cylindrical shell is composed of FG porous materials, reinforced with CNT-embedded nanocomposite layers to enhance its mechanical performance. The nonlinear governing motion equations are derived using Hamilton's principle. Mathematical solutions for the sandwich cylindrical shell are obtained through the application of Galerkin's and multiple-scale methods. The robustness of our results is substantiated by conducting a comprehensive comparison with prior research findings. The implications of our outcomes span multiple disciplines, including aerospace, automotive, civil engineering, mechanical engineering, and nanotechnology. These findings offer valuable insights into the vibrational behavior of such structures, which are of significant relevance in various real-world applications.

## 2. Model's description

In this section, the current model under examination is introduced. Fig. 1 displays a cylindrical shell comprising a sandwich structure, where an FG porous core is subjected to compression between layers composed of CNT-reinforced composite materials. The geometric components are depicted in the figure. For context, the entire assembly is assumed to be positioned on a Pasternak elastic foundation. Notations such as  $h_c$ ,  $h_t$ , and  $h_b$  are used to represent the thickness of the porous core, top face sheet, and bottom face sheet, respectively. The total thickness, denoted as  $h_{Total}$ , is calculated as the sum of  $h_t$ ,  $h_c$ , and  $h_b$ . The Cartesian coordinate system  $(x, y, z)$ , as shown in Fig. 1, is employed for this analysis. Furthermore, the model's geometric properties encompass dimensions such as length ( $a$ ), width ( $b$ ), and curvature radius ( $R$ ). It is important to note that throughout our analysis, a perfect bond is assumed between all three layers, indicating complete adhesion among them.

To enhance precision, the FSDT is employed for characterizing the displacement field within the sandwich model. According to this theory, it becomes feasible to represent the displacements at arbitrary points, whether situated within the FG porous core or the nanocomposite layers, in the following manner (Mindlin 1968):

$$\begin{aligned} u(x, y, z, t) &= u_0(x, y, t) + z\varphi_x(x, y, t) \\ v(x, y, z, t) &= v_0(x, y, t) + z\varphi_y(x, y, t) \\ w(x, y, z, t) &= w_0(x, y, t) \end{aligned} \quad (1)$$

In this context,  $u$ ,  $v$ , and  $w$  symbolize the displacements in the  $x$ -direction,  $y$ -direction, and the transverse displacement along the  $z$ -direction, respectively. Moreover,  $\theta_x$  and  $\theta_y$  refer to the rotations of the mid-plane around the  $y$  and  $x$  axes, respectively.

Employing von-Karman's assumptions, utilizing the previously mentioned displacement field, and accounting for nonlinearity, the strain field can be defined as follows (Mao and Zhang 2019):

$$\varepsilon_{xx} = \frac{\partial u}{\partial x} + \frac{1}{2} \left( \frac{\partial w_0}{\partial x} \right)^2 = \frac{\partial u_0}{\partial x} + z \frac{\partial \varphi_x}{\partial x} + \frac{1}{2} \left( \frac{\partial w_0}{\partial x} \right)^2, \quad (2)$$

$$\begin{aligned}\varepsilon_{yy} &= \frac{\partial v}{\partial y} + \frac{w}{R} + \frac{1}{2} \left( \frac{\partial w_0}{\partial y} \right)^2 \\ &= \frac{\partial v_0}{\partial y} + z \frac{\partial \varphi_y}{\partial y} + \frac{w_0}{R} + \frac{1}{2} \left( \frac{\partial w_0}{\partial y} \right)^2,\end{aligned}$$

$$\begin{aligned}\gamma_{xy} &= \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} + \frac{\partial w_0}{\partial x} \frac{\partial w_0}{\partial y} \\ &= \frac{\partial v_0}{\partial x} + z \frac{\partial \varphi_y}{\partial x} + \frac{\partial u_0}{\partial y} + z \frac{\partial \varphi_x}{\partial y} + \frac{\partial w_0}{\partial x} \frac{\partial w_0}{\partial y},\end{aligned}$$

$$\gamma_{xz} = \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} = \varphi_x + \frac{\partial w_0}{\partial x},$$

$$\gamma_{yz} = \frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} - \frac{v}{R} = \varphi_y + \frac{\partial w_0}{\partial y} - \frac{v_0}{R} - \frac{z}{R} \varphi_y$$

In this study, the precise position of the shell's neutral surface is determined and then employed in the subsequent relationships. This positioning is achieved to ensure satisfaction of the following equation (Radić 2018):

$$\int_{-\frac{h_{Total}}{2}}^{\frac{h_{Total}}{2}} c_{11}(z - C) dz = 0 \quad (3)$$

Here  $C$  denoting the distance of the neutral surface position from the mid-surface of the shell, can be determined by:

$$C = \frac{\int_{-\frac{h_{Total}}{2}}^{\frac{h_{Total}}{2}} c_{11} z dz}{\int_{-\frac{h_{Total}}{2}}^{\frac{h_{Total}}{2}} c_{11} dz} \quad (4)$$

The stress-strain relationships for both the FG porous core and CNT-reinforced composite layers are articulated as follows (Alhaifi *et al.* 2023, Kaveh *et al.* 2024):

$$\begin{bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \sigma_{yz} \\ \sigma_{xz} \\ \sigma_{xy} \end{bmatrix} = \begin{bmatrix} Q_{11} & Q_{12} & 0 & 0 & 0 \\ Q_{21} & Q_{22} & 0 & 0 & 0 \\ 0 & 0 & Q_{44} & 0 & 0 \\ 0 & 0 & 0 & Q_{55} & 0 \\ 0 & 0 & 0 & 0 & Q_{66} \end{bmatrix}^{c,f} \begin{bmatrix} \varepsilon_{xx} \\ \varepsilon_{yy} \\ \gamma_{yz} \\ \gamma_{xz} \\ \gamma_{xy} \end{bmatrix} \quad (5)$$

The elastic coefficients, denoted as  $c_{ij}$ , are specified individually for both the FG porous core and the CNT-reinforced composite layers, and the superscripts  $c$  and  $f$  denote the core and face sheets.

### 3. Nonlinear governing motion equations

The derivation of the nonlinear governing equations of motion utilizes Hamilton's principle, which takes into account the energies associated with strain and kinetic components as well as external work. This principle can be formally presented as follows (Arshid *et al.* 2024a, Jalaei and Civalek 2019):

$$\delta \int_{t_1}^{t_2} (U - K - V) dt = 0 \quad (6)$$

In this context,  $U$ ,  $T$ , and  $W$  represent the strain energy, kinetic energy, and external work, respectively. It is important to highlight that in this paper, the Pasternak foundation is assumed to bear the weight of the sandwich

model. This foundation has the unique capability to simulate both transverse and normal loads. The force exerted on the sandwich model by the foundation can be formally presented as follows (Arefi *et al.* 2018b, Mohammad-Rezaei Bidgoli and Arefi 2023c):

$$F_f = k_1 w_0 - k_2 \left( \frac{\partial^2 w_0}{\partial x^2} + \frac{\partial^2 w_0}{\partial y^2} \right) \quad (7)$$

The Winkler constant, symbolized as  $k_1$ , and the shear layer parameter, designated as  $k_2$ , characterize the characteristics of the Pasternak substrate. Consequently, the changes in external work induced by the presence of the Pasternak substrate can be formally articulated as follows (Arshid *et al.* 2022, 2024b):

$$\delta V = \int_A -F_f \left( 1 - \frac{h}{2R} \right) \delta w_0 dA \quad (8)$$

Next, the focus shifts to the calculation of strain energy and its variations. The strain energy of the model under examination is defined as follows (Zavari *et al.* 2024):

$$\delta U = \iiint_V (\sigma_{ij} \delta \varepsilon_{ij}) dV \quad (9)$$

Its variation is determined as follows:

$$\delta U = \iiint_V \left( \begin{aligned} &\sigma_{xx} \delta \varepsilon_{xx} + \sigma_{yy} \delta \varepsilon_{yy} + \sigma_{yz} \delta \gamma_{yz} \\ &+ \sigma_{xz} \delta \gamma_{xz} + \sigma_{xy} \delta \gamma_{xy} \end{aligned} \right) dV \quad (10)$$

By inserting the obtained strains into the equation mentioned above, the following expression is derived:

$$\delta U = \iiint_V \left[ \begin{aligned} &\sigma_{xx} \left( \frac{\partial \delta u_0}{\partial x} + z \frac{\partial \delta \varphi_x}{\partial x} + \frac{\partial w_0}{\partial x} \frac{\partial \delta w_0}{\partial x} \right) \\ &+ \sigma_{yy} \left( \frac{\partial \delta v_0}{\partial y} + z \frac{\partial \delta \varphi_y}{\partial y} + \frac{\delta w_0}{R} + \frac{\partial w_0}{\partial y} \frac{\partial \delta w_0}{\partial y} \right) \\ &+ \sigma_{yz} \left( \delta \varphi_y + \frac{\partial \delta w_0}{\partial y} - \frac{\delta v_0}{R} - \frac{z}{R} \delta \varphi_y \right) \\ &+ \sigma_{xz} \left( \delta \varphi_x + \frac{\partial \delta w_0}{\partial x} \right) \\ &+ \sigma_{xy} \left( \frac{\partial \delta v_0}{\partial x} + z \frac{\partial \delta \varphi_y}{\partial x} + \frac{\partial \delta u_0}{\partial y} + z \frac{\partial \delta \varphi_x}{\partial y} \right) \\ &+ \sigma_{xy} \left( \frac{\partial \delta w_0}{\partial x} \frac{\partial w_0}{\partial y} + \frac{\partial w_0}{\partial x} \frac{\partial \delta w_0}{\partial y} \right) \end{aligned} \right] dV \quad (11)$$

Through the separation of integrals with respect to  $z$  and the establishment of stress resultants, in addition to the application of variational formulations, we arrive at the following expression for the variations in strain energy:

$$\delta U = \iint_A \left[ \begin{aligned} &\left( -\frac{\partial H_{xx}}{\partial x} - \frac{\partial H_{xy}}{\partial y} \right) \delta u_0 \\ &+ \left( -\frac{\partial H_{yy}}{\partial y} - J_{yz} - \frac{\partial H_{xy}}{\partial x} \right) \delta v_0 \\ &+ \left( -\frac{\partial}{\partial x} \left( H_{xx} \frac{\partial w_0}{\partial x} \right) + J_{yy} \right) \\ &+ \left( -\frac{\partial}{\partial y} \left( H_{yy} \frac{\partial w_0}{\partial y} \right) - \frac{\partial H_{yz}}{\partial y} - \frac{\partial H_{xz}}{\partial x} \right) \delta w_0 \\ &+ \left( -\frac{\partial}{\partial x} \left( H_{xy} \frac{\partial w_0}{\partial y} \right) - \frac{\partial}{\partial y} \left( H_{xy} \frac{\partial w_0}{\partial x} \right) \right) \\ &+ \left( -\frac{\partial G_{xx}}{\partial x} + H_{xz} - \frac{\partial G_{xy}}{\partial y} \right) \delta \varphi_x \\ &+ \left( -\frac{\partial G_{yy}}{\partial y} + H_{yz} - G_{yz} - \frac{\partial G_{xy}}{\partial x} \right) \delta \varphi_y \end{aligned} \right] dA \quad (12)$$

in which:

$$(H_{xx}, G_{xx}) = \int_{-\frac{h_c}{2}-h_b-c}^{-\frac{h_c}{2}-c} \sigma_{xx}^b(1, z) dz + \int_{-\frac{h_c}{2}-c}^{+\frac{h_c}{2}-c} \sigma_{xx}^c(1, z) dz + \int_{+\frac{h_c}{2}-c}^{+\frac{h_c}{2}+h_t-c} \sigma_{xx}^t(1, z) dz, \quad (13)$$

$$(H_{yy}, G_{yy}, J_{yy}) = \int_{-\frac{h_c}{2}-h_b-c}^{-\frac{h_c}{2}-c} \sigma_{yy}^b\left(1, z, \frac{1}{R}\right) dz + \int_{-\frac{h_c}{2}-c}^{+\frac{h_c}{2}-c} \sigma_{yy}^c\left(1, z, \frac{1}{R}\right) dz + \int_{+\frac{h_c}{2}-c}^{+\frac{h_c}{2}+h_t-c} \sigma_{yy}^t\left(1, z, \frac{1}{R}\right) dz, \quad (14)$$

$$(H_{yz}, G_{yz}, J_{yz}) = \int_{-\frac{h_c}{2}-h_b-c}^{-\frac{h_c}{2}-c} k_s \sigma_{yz}^b\left(1, \frac{z}{R}, \frac{1}{R}\right) dz + \int_{-\frac{h_c}{2}-c}^{+\frac{h_c}{2}-c} k_s \sigma_{yz}^c\left(1, \frac{z}{R}, \frac{1}{R}\right) dz + \int_{+\frac{h_c}{2}-c}^{+\frac{h_c}{2}+h_t-c} k_s \sigma_{yz}^t\left(1, \frac{z}{R}, \frac{1}{R}\right) dz, \quad (15)$$

$$H_{xz} = \int_{-\frac{h_c}{2}-h_b-c}^{-\frac{h_c}{2}-c} k_s \sigma_{xz}^b dz + \int_{-\frac{h_c}{2}-c}^{+\frac{h_c}{2}-c} k_s \sigma_{xz}^c dz + \int_{+\frac{h_c}{2}-c}^{+\frac{h_c}{2}+h_t-c} k_s \sigma_{xz}^t dz, \quad (16)$$

$$(H_{xy}, G_{xy}) = \int_{-\frac{h_c}{2}-h_b-c}^{-\frac{h_c}{2}-c} \sigma_{xy}^b(1, z) dz + \int_{-\frac{h_c}{2}-c}^{+\frac{h_c}{2}-c} \sigma_{xy}^c(1, z) dz + \int_{+\frac{h_c}{2}-c}^{+\frac{h_c}{2}+h_t-c} \sigma_{xy}^t(1, z) dz \quad (17)$$

In the above relations,  $k_s$  is the shear correction factor of FSDT. On the other hand, in order to calculate the kinetic energy of the sandwich shell, the following formula can be applied (Arshid *et al.* 2024c, Kargar *et al.* 2021, Zhao *et al.* 2019):

$$\int_{t_1}^{t_2} K dt = \int_{t_1}^{t_2} \iiint_V \rho(z) \left(\frac{\partial u_i}{\partial t}\right)^2 dV dt \quad (18)$$

Regarding the defined displacements in Eq. (1), variations of the kinetic energy may be written as:

$$\int_{t_1}^{t_2} \delta K dt \quad (19)$$

$$= \int_{t_1}^{t_2} \iint_A \left[ \begin{aligned} & \left(-I_0 \frac{\partial^2 u_0}{\partial t^2} - I_1 \frac{\partial^2 \varphi_x}{\partial t^2}\right) \delta u_0 \\ & + \left(-I_0 \frac{\partial^2 v_0}{\partial t^2} - I_1 \frac{\partial^2 \varphi_y}{\partial t^2}\right) \delta v_0 \\ & + \left(-I_0 \frac{\partial^2 w_0}{\partial t^2}\right) \delta w_0 \\ & + \left(-I_1 \frac{\partial^2 u_0}{\partial t^2} - I_2 \frac{\partial^2 \varphi_x}{\partial t^2}\right) \delta \varphi_x \\ & + \left(-I_1 \frac{\partial^2 v_0}{\partial t^2} - I_2 \frac{\partial^2 \varphi_y}{\partial t^2}\right) \delta \varphi_y \end{aligned} \right] dAdt$$

where:

$$(I_0, I_1, I_2) = \int_{-\frac{h_c}{2}-h_b-c}^{-\frac{h_c}{2}-c} \rho^b(1, z, z^2) dz + \int_{-\frac{h_c}{2}-c}^{+\frac{h_c}{2}-c} \rho^c(1, z, z^2) dz + \int_{+\frac{h_c}{2}-c}^{+\frac{h_c}{2}+h_t-c} \rho^t(1, z, z^2) dz \quad (20)$$

In this case, applying Hamilton's principle and considering the obtained expressions for variations in strain and kinetic energies, as well as the inclusion of external work, results in the extraction of the following set of five nonlinear differential governing equations of motion:

$$\delta u_0: \frac{\partial H_{xx}}{\partial x} + \frac{\partial H_{xy}}{\partial y} = I_0 \frac{\partial^2 u_0}{\partial t^2} + I_1 \frac{\partial^2 \varphi_x}{\partial t^2}, \quad (21)$$

$$\delta v_0: \frac{\partial H_{yy}}{\partial y} + J_{yz} + \frac{\partial H_{xy}}{\partial x} = I_0 \frac{\partial^2 v_0}{\partial t^2} + I_1 \frac{\partial^2 \varphi_y}{\partial t^2}, \quad (22)$$

$$\delta w_0: \frac{\partial}{\partial x} \left( H_{xx} \frac{\partial w_0}{\partial x} \right) + \frac{\partial}{\partial y} \left( H_{yy} \frac{\partial w_0}{\partial y} \right) + \frac{\partial}{\partial x} \left( H_{xy} \frac{\partial w_0}{\partial y} \right) + \frac{\partial}{\partial y} \left( H_{xy} \frac{\partial w_0}{\partial x} \right) + \frac{\partial H_{yz}}{\partial y} + \frac{\partial H_{xz}}{\partial x} - J_{yz} - F_f \left( 1 - \frac{h}{2R} \right) = I_0 \frac{\partial^2 w_0}{\partial t^2}, \quad (23)$$

$$\delta \varphi_x: \frac{\partial G_{xx}}{\partial x} + \frac{\partial G_{xy}}{\partial y} - H_{xz} = I_1 \frac{\partial^2 u_0}{\partial t^2} + I_2 \frac{\partial^2 \varphi_x}{\partial t^2}, \quad (24)$$

$$\delta \varphi_y: \frac{\partial G_{yy}}{\partial y} + \frac{\partial G_{xy}}{\partial x} - H_{yz} + G_{yz} = I_1 \frac{\partial^2 v_0}{\partial t^2} + I_2 \frac{\partial^2 \varphi_y}{\partial t^2} \quad (25)$$

#### 4. Materials properties

Attention is now directed to the inner and outer layers, which are constructed using CNT-reinforced composite materials. These layers are composed of a polymer matrix with CNT fibers serving as the reinforcing phase. The elastic coefficients for these layers are defined as follows (Arshid *et al.* 2023c, Bidgoli and Arefi 2023):

$$Q_{11}^f = \frac{E_{11}}{1 - \nu_{12}\nu_{21}}, \quad Q_{22}^f = \frac{E_{22}}{1 - \nu_{12}\nu_{21}}, \quad (26)$$

$$Q_{12}^f = \nu_{21} Q_{11}^f, \quad Q_{44}^f = G_{23},$$

$$Q_{55}^f = G_{13}, \quad Q_{66}^f = G_{12}$$

Here, it's important to highlight that the knowledge of the mechanical properties of the composite layer is imperative for the mathematical derivation and solution of the governing equations. To achieve this, the extended rule of mixture (ERM) is employed to combine the mechanical properties of the matrix and reinforcing phases, thereby establishing the mechanical properties of the entire composite layer. In accordance with the ERM approach, the material properties of the nanocomposite layers are as follows (Allahkarami *et al.* 2018, Arshid *et al.* 2023d):

$$\nu_{12} = V_{main}^{CNT} \nu_{12}^{CNT} + V^P \nu^P, \quad (27)$$

$$\rho_c = V^{CNT} \rho^{CNT} + V^P \rho^P, \quad (28)$$

$$E_{11} = \eta_1 V^{CNT} E_{11}^{CNT} + V^P E^P, \quad (29)$$

$$\frac{\eta_3}{G_{12}} = \frac{V^{CNT}}{G_{12}^{CNT}} + \frac{V^P}{G^P}, \quad \frac{\eta_2}{E_{22}} = \frac{V^{CNT}}{E_{22}^{CNT}} + \frac{V^P}{E^P} \quad (30)$$

In these equations, the letter  $P$  signifies the polymeric matrix, and  $\eta_k$  (where  $k$  takes on values of 1, 2, and 3) represents the CNT efficiency parameter. It's important to note that  $\eta_k$  depends on the total volume of CNTs present in the composite material, as detailed in Table 1 (Jeyaraj and Rajkumar 2013).

$V^P$  represents the volume fraction for the matrix, while  $V^{CNT}$  stands for the volume fraction pertaining to the CNTs. It's essential to emphasize that the sum of  $V^P$  and  $V^{CNT}$  must equal 1, indicating the entirety of the material composition.

Furthermore, it's crucial to distinguish  $V^{CNT}$  from  $V_{main}^{CNT}$  as  $V^{CNT}$  is dependent on the dispersion pattern of CNTs within the polymer matrix phase. In other words,  $V^{CNT}$  varies based on the CNT dispersion pattern. The expressions for  $V^{CNT}$  and  $V_{main}^{CNT}$  can be formulated as follows (Amir *et al.* 2020a):

$$V_{main}^{CNT} = \frac{w^{CNT}}{w^{CNT} + (\rho^{CNT}/\rho^P) - (\rho^{CNT}/\rho^P)w^{CNT}} \quad (31)$$

$$V_i^{CNT} = \begin{cases} V_{main}^{CNT} & FG-UU \\ \left[1 - \frac{2}{h_i} \left(z \pm \frac{h_c + h_i}{2}\right)\right] V_{main}^{CNT} & FG-AV \\ \left[1 + \frac{2}{h_i} \left(z \pm \frac{h_c + h_i}{2}\right)\right] V_{main}^{CNT} & FG-VA \\ 2 \left[1 - \frac{2}{h_i} \left(z \pm \frac{h_c + h_i}{2}\right)\right] V_{main}^{CNT} & FG-OO \\ \left[\frac{4}{h_i} \left(\left|z \pm \frac{h_c + h_i}{2}\right|\right)\right] V_{main}^{CNT} & FG-XX \end{cases} \quad i = t, b \quad (32)$$

In the equations mentioned above, the subscripts 't' and 'b' are abbreviations for 'Top' and 'Bottom,' respectively, arranged sequentially. Additionally,  $w$  signifies the mass fraction, while FG-AV, FG-VA, FG-OO, and FG-XX indicate the dispersion patterns of CNTs within the matrix phase. Specifically, 'UU' represents the uniform dispersion pattern.

Table 1 Efficiency parameters of CNTs (Jeyaraj and Rajkumar 2013)

| $V_{main}^{CNT}$ | $\eta_1$ | $\eta_2$ | $\eta_3$ |
|------------------|----------|----------|----------|
| 0.12             | 0.137    | 1.022    | 0.715    |
| 0.17             | 0.142    | 1.626    | 1.138    |
| 0.28             | 0.141    | 1.585    | 1.109    |

Now, let's delve into the definitions of elastic coefficients for the FG porous core. As highlighted in the previous section, a porous layer serves as the core of the aforementioned sandwich cylindrical shell. Consequently, the elastic coefficients are assumed to have the following values (Amir *et al.* 2020b, Khoddami Maraghi *et al.* 2022):

$$Q_{11}^c = Q_{22}^c = \frac{E_c(z)}{1 - \nu_c^2},$$

$$Q_{12}^c = \frac{\nu_c E_c(z)}{1 - \nu_c^2},$$

$$Q_{ii}^c = G_c(z), \quad i = 4, 5, 6 \quad (33)$$

in which  $E_c(z)$ ,  $G_c(z)$ , and  $\nu_c$  indicate elastic modulus, shear modulus and Poisson's ratio of the FG porous core, respectively.

The elastic constants of the core shell are influenced by the type of porosity distribution within it, leading to variations in both the elastic modulus and core density. Among the three established types of porosity distribution (uniform, non-uniform, and even), this paper focuses on the non-uniform distribution, specifically termed "Distribution 1". Under this condition, the following relationships define the elastic constants (Qin *et al.* 2020):

$$E_c(z) = E_0 \left[1 - e_0 \cos \left(\left\{\frac{\pi}{2h_c} \left(z + \frac{h_c}{2}\right)\right\}\right)\right], \quad (34)$$

$$\rho_c(z) = \rho_0 \left[1 - e_m \cos \left(\left\{\frac{\pi}{2h_c} \left(z + \frac{h_c}{2}\right)\right\}\right)\right], \quad (35)$$

$$e_m = 1 - \sqrt{1 - e_0} \quad (36)$$

These relationships are expressed in terms of  $e_0$  which represents the porosity coefficient, and  $e_m$  which is the mass density factor. In the "Distribution 2" pattern, it's noteworthy that there is an independence of thickness parameters in the modulus and density equations, which is associated with uniform porosity distribution (Hou *et al.* 2023):

$$E_c(z) = E_0 [1 - e_0 \zeta], \quad (37)$$

$$\rho_c(z) = \rho_0 \sqrt{1 - e_0 \zeta}, \quad (38)$$

$$\zeta = \frac{1}{e_0} - \frac{1}{e_0} \left(\frac{2}{\pi} \sqrt{1 - e_0} - \frac{2}{\pi} + 1\right)^2 \quad (39)$$

In the scenario of even porosity distribution, referred to as "Distribution 3," pores are symmetrically distributed across the porous layer. The corresponding distribution of elastic modulus and density can be defined as follows:

Table 2 Material properties of the CNT-reinforced composite layers (Alibeigloo 2013)

| Properties                  | SWCNTs | PVDF             |
|-----------------------------|--------|------------------|
| $\nu$                       | 0.175  | 0.34             |
| $\rho$ (kg/m <sup>3</sup> ) | 1400   | 1150             |
| $E_{11}$ (TPa)              | 5.6466 | $E_m$ (GPa)= 2.5 |
| $E_{22}=E_{33}$ (TPa)       | 7.0800 |                  |
| $G_{12}=G_{23}=G_{13}$      | 1.9445 |                  |

Table 3 Validation the results in simpler case with previously published work.

| R/a | Source                     | k      |        |        |        |        |
|-----|----------------------------|--------|--------|--------|--------|--------|
|     |                            | 0.5    | 1      | 2      | 5      | 10     |
| 5   | Kiani <i>et al.</i> (2012) | 7.0497 | 6.7171 | 6.4280 | 6.1525 | 5.9721 |
|     | Present                    | 7.0716 | 6.7387 | 6.4354 | 6.1720 | 5.9773 |
| 10  | Kiani <i>et al.</i> (2012) | 5.8314 | 5.5586 | 5.3502 | 5.1711 | 5.0285 |
|     | Present                    | 5.8492 | 5.5729 | 5.3624 | 5.1843 | 5.0494 |
| 20  | Kiani <i>et al.</i> (2012) | 5.4839 | 5.2282 | 5.0442 | 4.8944 | 4.7629 |
|     | Present                    | 5.4936 | 5.2376 | 5.0527 | 4.9018 | 4.7748 |

$$E_c(z) = E_0 \left[ 1 - e_0 \cos \left( \frac{\pi z}{h_c} \right) \right], \quad (40)$$

$$\rho_c(z) = \rho_0 \left[ 1 - e_m \cos \left( \frac{\pi z}{h_c} \right) \right] \quad (41)$$

In the aforementioned equations, it's important to note that  $E_0$  represents the maximum elastic modulus, while  $\rho_0$  stands for the maximum density within the FG porous core.

## 5. Semi-analytical solution process

The nonlinear governing equations of motion under simply supported boundary conditions are effectively solved through the application of Galerkin's method in conjunction with the multiple scale method. In this context, the following functions are defined (Mohammad-Rezaei Bidgoli and Arefi 2023a):

$$\begin{aligned} u_0(x, y, t) &= \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \tilde{U}_{mn}(t) \cos(\alpha x) \sin(\beta y), \\ v_0(x, y, t) &= \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \tilde{V}_{mn}(t) \sin(\alpha x) \cos(\beta y), \\ w_0(x, y, t) &= \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \tilde{W}_{mn}(t) \sin(\alpha x) \sin(\beta y), \\ \varphi_x(x, y, t) &= \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \tilde{N}_{mn}(t) \cos(\alpha x) \sin(\beta y), \\ \varphi_y(x, y, t) &= \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \tilde{D}_{mn}(t) \sin(\alpha x) \cos(\beta y) \end{aligned} \quad (42)$$

Here, where  $\beta$  is defined as  $n\pi/b$ , and  $\alpha$  is represented as  $m\pi/a$ . The coefficients  $\tilde{U}$ ,  $\tilde{V}$ ,  $\tilde{W}$ ,  $\tilde{N}$ , and  $\tilde{D}$  represent the unknowns in this context. The values  $(m, n)$  correspond to the mode numbers along the  $x$  and  $y$  directions. Upon substituting these solutions into the governing equations of motion and subsequently applying Galerkin's method (Dzung *et al.* 2024, Ha *et al.* 2024), taking certain assumptions into account, the final governing equation of motion in the time domain is obtained as follow:

$$\frac{d^2 \tilde{W}_{mn}(t)}{dt^2} + \alpha_1 \tilde{W}_{mn}(t) + \alpha_2 \tilde{W}_{mn}^2(t) + \alpha_3 \tilde{W}_{mn}^3(t) = 0 \quad (43)$$

Also, the initial conditions are as below:

$$\tilde{W}_{mn}(t=0) = \frac{W_{Max}}{h} = A, \quad \frac{d\tilde{W}_{mn}(t)}{dt}(t=0) = 0 \quad (44)$$

The explicit solution of the nonlinear equations leads to following result (Babaei and Eslami 2021):

$$\frac{\omega_{NL}}{\omega_L} = \left[ 1 + \frac{9\alpha_1\alpha_3 - 10\alpha_2^2}{12\alpha_1^2} A^2 \right]^{1/2} \quad (45)$$

## 6. Debate on the results

This section is dedicated to presenting the results, aiming to offer a comprehensive understanding of the nonlinear vibrational behavior exhibited by the under-evaluation model. The FG porous core has  $E_c=200 \times 10^9$  Pa,  $\rho_c=7850$  kg/m<sup>3</sup>, and  $\nu_c=0.33$ . Additionally, Table 2 provides the properties of the matrix (Epoxy) and the reinforcement (CNTs), serving as constituents of the inner and outer layers (Alibeigloo 2013).

To ensure the code's validity and enhance its reliability, results concerning the impacts of geometrical and material properties on linear and normalized natural frequencies were obtained using the current code. These results were subsequently compared with those reported by Kiani *et al.* (2012) in Table 3. This Table presents the results for  $a/b=1$ ,  $a/h=50$ , and various values of the  $R/a$  ratio and the material gradient index ( $k$ ) of FG materials. Upon examination of Table 3, it becomes apparent that the results from the mentioned paper (Kiani *et al.* 2012), while displaying minor differences, can be attributed to different assumptions or the utilization of varying displacement field formulations. Consequently, the code presented in the current paper can be considered reliable, and its results can offer dependable insights into the vibrational behavior of such lightweight structures.

Now it turns to consider the current work's results. The Fig. 2 is presented to study the effect of porosity coefficient and also,  $W_{Max}/h_{Total}$  ratio on the nonlinear-to-linear frequency ratio. It is seen that by increasing the porosity coefficient, the results tend to reduce. Increasing porosity typically reduces the effective mass of the shell because voids within the material have less mass compared to the solid material. In structural dynamics, the frequency of a structure is inversely proportional to the square root of its mass. Therefore, a reduction in mass due to increased porosity can lead to a decrease in the frequency. Porous

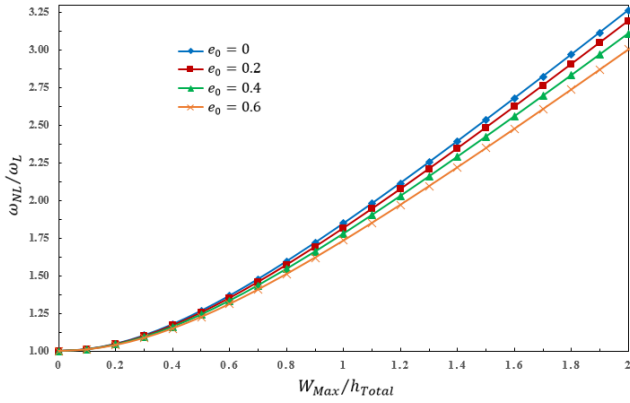


Fig. 2 Effect of porosity coefficient and  $W_{Max}/h_{Total}$  ratio on the results

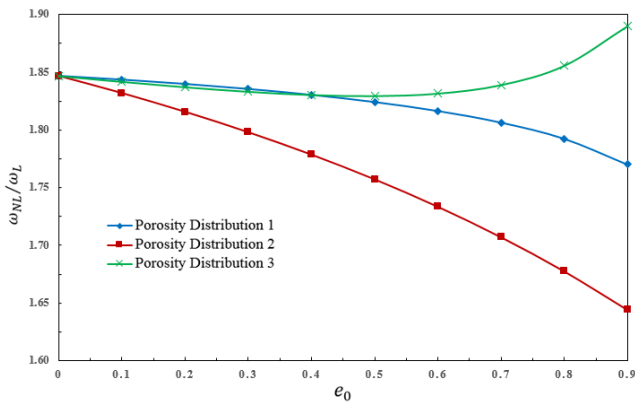


Fig. 3 Impact of porosity coefficient and different types of pores distribution on the results

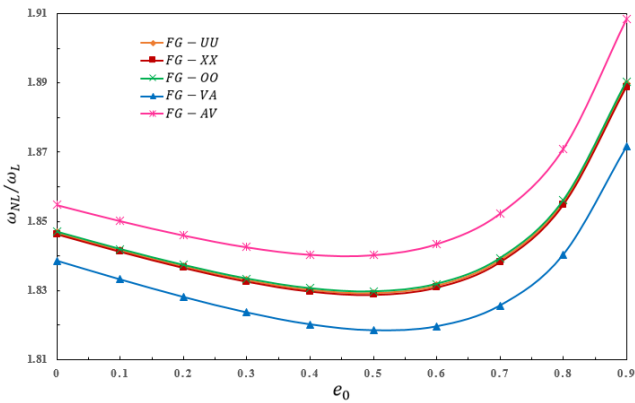


Fig. 4 Influence of porosity coefficient and different pattern of CNTs' dispersion in the inner and outer layer on the frequencies

materials are often less stiff than their solid counterparts. This reduced stiffness can also contribute to a decrease in frequency because stiffness is directly related to the square of the frequency. It's important to note that the exact relationship between porosity and frequency depends on the specific geometry and material properties of the shell, as well as the boundary conditions and the type of vibrations or oscillations under consideration. In some cases, increased porosity may lead to a significant reduction in frequency,

Table 4 Exact values of the results for CNTs' different distribution patterns and  $W_{Max}/h_{Total}$  ratio

| $W_{max}$<br>$h_{Total}$ | CNTs Distribution Pattern |        |        |        |        |
|--------------------------|---------------------------|--------|--------|--------|--------|
|                          | FG-UU                     | FG-XX  | FG-OO  | FG-VA  | FG-AV  |
| 0                        | 1.0000                    | 1.0000 | 1.0000 | 1.0000 | 1.0000 |
| 0.1                      | 1.0107                    | 1.0107 | 1.0107 | 1.0104 | 1.0110 |
| 0.2                      | 1.0422                    | 1.0421 | 1.0422 | 1.0410 | 1.0434 |
| 0.3                      | 1.0927                    | 1.0926 | 1.0928 | 1.0902 | 1.0953 |
| 0.4                      | 1.1597                    | 1.1594 | 1.1598 | 1.1555 | 1.1640 |
| 0.5                      | 1.2405                    | 1.2401 | 1.2407 | 1.2344 | 1.2468 |
| 0.6                      | 1.3327                    | 1.3322 | 1.3329 | 1.3244 | 1.3410 |
| 0.7                      | 1.4340                    | 1.4333 | 1.4342 | 1.4235 | 1.4445 |
| 0.8                      | 1.5426                    | 1.5418 | 1.5429 | 1.5299 | 1.5554 |
| 0.9                      | 1.6572                    | 1.6562 | 1.6575 | 1.6421 | 1.6722 |
| 1                        | 1.7764                    | 1.7753 | 1.7768 | 1.7591 | 1.7938 |
| 1.1                      | 1.8996                    | 1.8983 | 1.9000 | 1.8800 | 1.9192 |
| 1.2                      | 2.0259                    | 2.0245 | 2.0263 | 2.0040 | 2.0478 |
| 1.3                      | 2.1548                    | 2.1533 | 2.1553 | 2.1307 | 2.1790 |
| 1.4                      | 2.2859                    | 2.2842 | 2.2864 | 2.2595 | 2.3123 |
| 1.5                      | 2.4188                    | 2.4170 | 2.4193 | 2.3901 | 2.4474 |
| 1.6                      | 2.5532                    | 2.5512 | 2.5538 | 2.5223 | 2.5840 |
| 1.7                      | 2.6889                    | 2.6868 | 2.6895 | 2.6558 | 2.7220 |
| 1.8                      | 2.8257                    | 2.8235 | 2.8264 | 2.7904 | 2.8610 |
| 1.9                      | 2.9635                    | 2.9611 | 2.9642 | 2.9260 | 3.0010 |
| 2                        | 3.1021                    | 3.0996 | 3.1029 | 3.0624 | 3.1418 |

while in others, the effect may be less pronounced. Also, it is observed that by increasing the  $W_{Max}/h_{Total}$  ratio, the results enhance significantly. The following numerical values and specifications are used to extract this figure:

$h_c = 0.005m, h_b = h_t = 0.001m, a = b = 20h_{Total}, R = 10a, V_{main}^{CNT} = 0.28, e_0 = 0.3, k_1 = k_2 = 0$ , and Distribution 3 of porosity.

Curvatures plotted in Fig. 3 are responsible to examine the impacts of both porosity distribution types and also, coefficient. It is understood that increasing porosity in two types, i.e., Porosity distribution 1 and 2 leads to reduction in the results, while for Porosity distribution 3 that is related to the symmetric pattern, the results enhance a bit. It occurs due to reduction or enhancement in stiffness-to-mass ratio of the shell. The following numerical values and specifications are used to extract this figure:

$$h_c = 0.005m, h_b = h_t = 0.001m, a = b = 20h_{Total}, \\ R = 10a, V_{main}^{CNT} = 0.12, W_{max}/h_{Total} = 1, \\ FG - UU, k_1 = k_2 = 0$$

Results plotted in Fig. 4 can provide a useful physical vision toward different patterns of dispersion of CNT in the inner and outer layers. It is found that FG-AV and FG-VA patterns lead to the most and the least values of the results and the other patterns outcomes are too close to each other. The following numerical values and specifications are used to extract this figure:

$$h_c = 0.005m, h_b = h_t = 0.001m, \\ a = b = 20h_{Total}, R = 10a, V_{main}^{CNT} = 0.12,$$

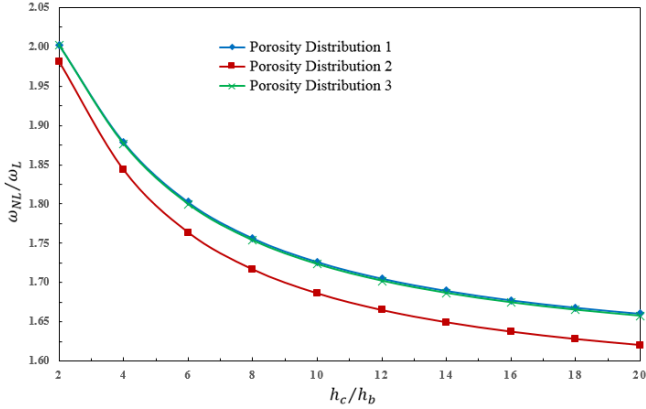


Fig. 5 Effect of layers' thicknesses ratio on the results

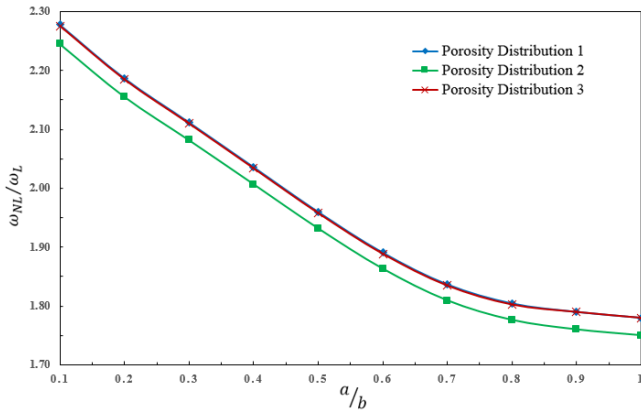


Fig. 6 Effect of geometrical parameters of the shell and porosity distribution on the results

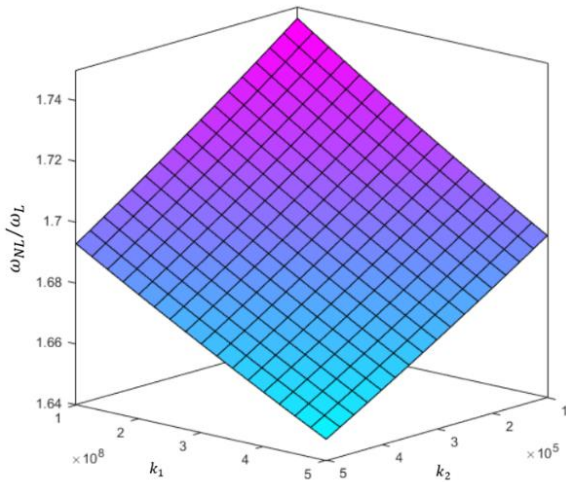


Fig. 7 Impact of Pasternak foundation's parameters on the results

$$W_{max}/h_{Total} = 1, Porosity\ Distribution\ 3, k_1 = k_2 = 0$$

Exact values of the results toward different patterns of CNTs' distribution and also,  $W_{Max}/h_{Total}$  ratio are listed in Table 4. This table confirms the previous findings about the effect of CNTs' distribution patterns on the results. The following numerical values and specifications are used to extract this table:

$$h_c = 0.005m, h_b = h_t = 0.001m,$$

$$a = b = 20h_{Total}, R = 10a, V_{main}^{CNT} = 0.28, \\ Porosity\ Distribution\ 3, e_0 = 0.3, k_1 = k_2 = 0$$

Fig. 5 shows the effect of layers' thicknesses ratio on the results. Regarding this figure it can be understood that increasing the core-to-faces thickness ratio leads the results to a reduction. In fact, the stiffness of the inner and outer layers is much more than the FG porous core and by keeping the total thickness constant, increasing the core's thickness means a reduction in inner and outer layers thickness which leads to reduction in stiffness of the whole structure. Therefore, the results are reduced, too. It occurs for all three-porosity distribution pattern that are investigated in this work. The following numerical values and specifications are used to extract this figure:

$$a = b = 20h_{Total}, R = 10a, V_{main}^{CNT} = 0.12, \\ W_{max}/h_{Total} = 1, FG - UU, k_1 = k_2 = 0, e_0 = 0.3$$

In Fig. 6, the influence of geometrical parameter of the shell, i.e., its length-to-width ratio is considered and it is seen due to a significant reduction in the shell's stiffness, consequently, the results are reduced. The following numerical values and specifications are used to extract this figure:

$$h_c = 0.005m, h_b = h_t = 0.001m, R = 10a, \\ V_{main}^{CNT} = 0.17, W_{max}/h_{Total} = 1, FG - UU, \\ k_1 = k_2 = 0, e_0 = 0.3, Area = Constant$$

And finally in Fig. 7, the effect of Pasternak foundation parameters is observed. As can be expected, increasing both these two parameters, i.e., springs and shear layer parameters, the rigidity of the structure enhance and accordingly, the results are also enhanced. The following numerical values and specifications are used to extract this figure:

$$h_c = 0.005m, h_b = h_t = 0.001m, a = b = 20h_{Total}, \\ R = 10a, V_{main}^{CNT} = 0.28, Porosity\ Distribution\ 3, \\ e_0 = 0.3, FG - UU, W_{max}/h_{Total} = 1$$

## 7. Conclusions

The nonlinear vibrational characteristics of sandwich cylindrical shells are considered in this study. To bring our findings closer to practical applications and align them with experimental studies, the FSDT is employed to model various points within the structure. The core of these cylindrical shells is comprised of FG porous materials, which are further reinforced with layers of nanocomposites embedded with CNTs to enhance their mechanical performance. The governing equations for nonlinear motion are derived using Hamilton's principle. Subsequently, mathematical solutions for these sandwich cylindrical shells are obtained through the application of Galerkin's and multiple-scale methods. A comprehensive comparison with prior research findings is conducted to substantiate the robustness of our results. Lastly, the influence of different parameters on the results is considered. It is generally observed that the results tend to be reduced by an increase in the porosity coefficient. However, in cases of a symmetric distribution of porosity, a slight enhancement in the results is noted. Moreover, an increase in the  $W_{Max}/h_{Total}$

ratio leads to a significant enhancement in the results. The most and least values of the results are found to be associated with different porosity patterns, namely FG-AV and FG-VA, respectively, while the outcomes for other patterns are observed to be very close to each other. An increase in the core-to-faces thickness ratio is seen to result in a reduction in the outcomes. Additionally, due to a significant reduction in the shell's stiffness, the results are decreased. Furthermore, an enhancement in the rigidity of the structure is observed by increasing both foundation parameters, namely springs and shear layer parameters, leading to an improvement in the results.

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