

Analytical nonlocal elasticity solution and ANN approximate for free vibration response of layered carbon nanotube reinforced composite beams

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Abstract. This article investigates the free vibration behavior of carbon nanotube reinforced composite (CNTRC) beams embedded using variational analytical methods and artificial neural networks (ANN). The material properties of layered functionally graded CNTRC (FG-CNTRC) beams are estimated using nonlocal parameters modified power-law with different types of CNT distributions through the thickness direction of the beam. Adopting Eringen's nonlocal elasticity theory to capture the small size effects, the nonlocal governing equations are derived and solved using the analytical method. And also, the problem was analyzed using the ANN method. The architecture of the proposed ANN model is 3-9-1. In the experiments, we used 112 different data to predict the natural frequency using ANN. Based on the nonlocal differential constitutive relations of Eringen, the equations of motion as well as the boundary conditions of the beam are derived using Hamilton's principle. The classical beam theory is used to formulate a governing equation for predicting the free vibration of laminated CNTRC beams. According to the experimental results, the prediction ability of the ANN model is very good and the natural frequency can be predicted in ANN without attempting any experiments.

Keywords: artificial neural networks; carbon nanotube; composite beam; free vibration; nonlocal elasticity theory

1. Introduction

The carbon nanotubes (CNTs) have become utilized in many essential fields including engineering and have received attentions amid scientists. The application of CNT reinforcement structures in many industries is rising due to their excellent properties such as high strength, low weight, resistance to fatigue and are useable in continuous yarns (Jam and Kiani 2015, Khosravi *et al.* 2019a, b, Khaniki and Ghayesh 2020, Babaei *et al.* 2021, Madenci *et al.* 2023). The single-walled CNT has a cylindrical shape that is continuously extruded with a single sheet of graphene about 1 nm in diameter and centimeter long (Bourada *et al.* 2020). In contrast, multi-walled CNT has an array of such cylinders that are concentrically formed and separated by 0.35 nm, similar to basal plane separation in graphite with diameters of 2 nm to 100 nm and lengths of tens of microns (Asghar *et al.* 2020). In traditional methods, nanotubes are uniformly or randomly distributed, therefore, the improvement of their special mechanical properties may not be perfect (Qin *et al.* 2020). Nanotechnology and nanoscience are attractive research topics that have been taken into considerations in the past decades for this purpose (Kannan *et al.* 2012, Mirzaei and Kiani 2015, Khater and Abd el Gawaad 2016, Kumar 2018, Eltahir *et al.* 2019, Asghar *et al.* 2020, Hosseini and Khosravi 2020, Hussain *et al.* 2020a,

b, c, Farazin and Mohammadimehr 2021).

Carbon nanotube reinforced composites (CNTRCs) are a new generation of composite materials in which the ratio of the two in-plane Young's modulus is more than 40 and the in-plane shear modulus is very small compared to the longitudinal Young's modulus (Shen *et al.* 2020). Therefore, the auxiliary CNTRC laminate is expected and the magnitude of the Poisson ratio can be increased in some directions. CNTRC laminated structures have potential use in the industry and engineering fields such as mechanical, civil, aerospace as load bearing instrument for ships, building and aircrafts, etc (Kiani 2016).

Methods such as atomistic modeling and continuum mechanics are used to understand the mechanical performance of CNTs (Eltahir *et al.* 2019). Eberhardt and Wallmersperger (2015) studied a modified molecular structural mechanics model to determine the elastic properties of CNTs. Ghadyani *et al.* (2016) proposed an algorithm based on a numerical finite element analysis, to relate the physical geometry to the elastic properties of asymmetric SWCNTs. Kumar (2018) presented mechanical vibration of double-walled carbon nanotubes with inter-tube Van der Waals forces.

In the literature, plenty of theories to analyze the small-scale effect, including surface effect, modified coupled stress method, nonlocal elasticity theory and nonlocal strain gradient theory. The small-scale effect is very important in nanoscale structures, so several dimension-dependent continuum models have been developed. However, the most used continuity theory to investigate small-scale structures

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is the nonlocal elasticity model proposed by Eringen (1983). The application of nonlocal elasticity in micro and nanomaterials has received much attention among the nanotechnology community recently. The nonlocal elasticity model is widely used in small-scale materials for bending, vibration and buckling analyzes of micro/nano structures (Bounouara *et al.* 2016, Mehar *et al.* 2016, Mouffoki *et al.* 2017, Karami *et al.* 2018, Youcef *et al.* 2018). Aydogdu (2012) applied nonlocal elasticity theory in axial vibration analysis of single-walled carbon nanotubes embedded in elastic medium. Şimşek and Yurtcu (2013) presented an analytical solution for bending and buckling of FG nanobeams using the nonlocal Timoshenko beam theory. Free vibration analysis of double walled carbon nanotubes (DWCNT) embedded in elastic medium with initial imperfection is investigated by Ehyaei and Daman (2017). Timoshenko beam theory is employed. However the nonlocal theory is used for modeling the nano scale of nanotube. In addition, the governing equations were obtained due to the Hamilton principle. Khosravi and Hosseini (2020) conducted the nonlocal model to investigate the small-scale torsional behavior of the single-walled CNTs for free case and for the state in which model is subjected to the linear and harmonic torques. Static bending and buckling of the FG porous nanoscale beam are investigated based on the nonlocal Timoshenko beam theory including neutral surface position by Bensaid and Guenanou (2017). Hosseini *et al.* (2020) investigated axial vibration of FG nanobeams using nonlocal elasticity theory under clamped-clamped and clamped-free boundary conditions. Power law, exponential law and sigmoid law were applied as grading laws to examine the effect of the material distribution on axial vibration of the FG nanobeams. Eltaher *et al.* (2016) investigated the effects of thermal load and shear force on the buckling of nanobeams. The nonlocal Eringen elasticity model was adopted to account a size-dependence at the nano-scale. A simple and refined nonlocal hyperbolic higher-order beam theory is proposed for bending and vibration response of nanoscale beams by Kheroubi *et al.* (2016). The large amplitude free vibrations of a sandwich beam with stiff core and CNTRC face sheets are analyzed by Mirzaei and Kiani (2016). The equivalent single-layer theory of Timoshenko is used to construct the Hamiltonian of the beam under the von Kármán type of geometrical nonlinearity assumptions.

Application of classical continuum theory to the analysis of nanostructures is inappropriate because classical theories lack accountability for size effects resulting from small scale. In classical elasticity theory, the stress at a point depends only on the strain at the same point, whereas in the nonlocal elasticity theory the stress at a point is a function of the strains at all points of the continuum (Wang and Li 2014). In this way, the nonlocal continuum theory contains information about the long-range forces between atoms, and the internal length scale is simply included as a material parameter in the constitutive equations to capture the small-scale effect. The propagation characteristics of the elastic wave in fluid-conveying double-walled carbon nanotubes are investigated by Wang *et al.* (2010). The nonlocal continuum theory is applied to discuss the dispersion

relation with small scale effects. The effects of the viscous fluid on the propagation characteristics of elastic waves in carbon nanotubes are studied Wang *et al.* (2011). Based on the nonlocal continuum theory, the small scales effects are also considered. Numerical simulations are performed with the consideration of different scale coefficients to discuss the influence of the viscous fluid. Thai and Vo (2012) proposed a nonlocal shear deformation beam theory for bending, buckling and vibration of homogeneous nanobeams based on Eringen's nonlocal elasticity theory. The effects of the axial load and the elastic matrix on the flexural wave in the carbon nanotube are studied by Wang *et al.* (2012). Based on the nonlocal continuum theory and the Timoshenko beam model, the equation of the flexural wave motion is derived. The dispersion relation between the frequency and the wave number is illustrated. Ebrahimi and Barati (2016) investigated the buckling behavior of shear deformable piezoelectric nanoscale beams made of functionally graded materials embedded in Winkler-Pasternak elastic medium and subjected to an electromagnetic field. To consider the small size effects, Eringen's nonlocal elasticity theory was adopted. This article investigates the vibration behavior of magneto-electro-elastic functionally graded nanobeams embedded in a two-parameter elastic foundation using a third-order parabolic shear deformation beam theory. Based on Eringen's nonlocal elasticity theory which captures the small size effects and using the Hamilton's principle, the nonlocal governing equations of motions were derived and then solved analytically. A new nonlocal higher order shear deformation theory was developed for buckling properties of a single graphene sheet by Bouadi *et al.* (2018). The proposed nonlocal higher order shear deformation theory contains a new displacement field which incorporates undetermined integral terms and contains only two variables. The length scale parameter was considered in the present formulation by employing the nonlocal differential constitutive relations of Eringen. Mokhtar *et al.* (2018) presented a novel shear deformation theory for buckling analysis of single-layer graphene sheet based on nonlocal elasticity theory.

Experimental simulations on nanostructures are difficult, so some researchers studied and reported theoretical analyses of CNTRC structures using classical theories such as Euler-Bernoulli, Kirchhoff etc. and first order/high order shear deformation theories (Shekari *et al.* 2017, Chavan and Lal 2018, Ebrahimi and Farazmandnia 2018). Madenci (2021) presented efficient energy functional for CNTRC nanobeams based on high order shear deformation theory via mixed finite element method. Madenci and Gülcü (2020) proposed a model based on artificial neural networks was performed well for the maximum displacement of FGM beams such as Young modulus ratio, length to thickness ratio, the effect of the shear correction factor, effect of the power-law exponent. Ebrahimi and Barati (2018) presented an analytical solution of FG piezoelectric nanobeams by using a developed third-order shear deformation theory. Duc *et al.* (2018) investigated dynamic analysis of imperfect rectangular nanocomposite multilayer organic solar cells using the classical plate theory. Wu *et al.* (2015)

examined the free vibration and buckling behavior of sandwich beams reinforced with FG-CNTRCs face sheets based on Timoshenko beam theory. A new nonlocal higher order shear deformation theory is developed for buckling properties of a single graphene sheet by Bouadi *et al.* (2018). A new nonlocal higher order shear deformation theory was developed for buckling properties of a single graphene sheet by Bouadi *et al.* (2018). The proposed nonlocal higher order theory contains a new displacement field that incorporates undetermined integral terms and contains only two variables. An important and important issue in fiber reinforced polymer (FRP) composites is the reinforcement of the interlayer region. The properties of the regions between weak layers limit the overall performance of laminated FRP composite structures. In FGMs, the change of material properties has a functional transition. So there is no delamination problem (Uyaner and Yar 2019, Vedernikov *et al.* 2020, Alajarmeh *et al.* 2021, Korotkov *et al.* 2021, Madenci 2023).

Recently soft computing techniques such as artificial neural networks (ANNs), genetic algorithm and fuzzy logic have been used extensively for damage assessment with varied success (Hakim and Razak 2013). ANNs, inspired by the human brain, are powerful tools used to solve many real-life problems, which are considered a powerful method in the field of structural dynamics. ANNs are usually employed when the relationship between the input and output is complicated or if the application of another available method takes a large amount of computational time and if the effort is very expensive. Gülcü (2021) developed an ANN model to predict the compressive strength of the rubberized concrete. On the other hand, there are not enough applications of ANNs on the analysis of nanostructures in the literature.

In this study, analytical solutions for free vibration of CNTRC nanobeams with arbitrary boundary conditions considering the nonlocal parameter and ANN approximate are presented. The material properties of FG- CNTRCs are assumed to be graded through the thickness direction according to several distributions namely UD, FG-X, FG-V and FG-O. This distribution forms are given in Table 1. In ANNs approach, the ANN model was created by taking into account the natural frequency of the program and various key parameters that may affect its testing. The use of ANN modeling to analyze the free vibration of layered FG-CNTRC beams was validated by performing comprehensive convergence studies. The analytical and ANNs solutions for vibration analysis demonstrated the effects of nonlocal parameter, CNT distribution, layer angle and boundary conditions.

2. Formulation

The optimal design of beams has an important role in the development of nanotechnology. The mechanical parameters including maximum deflection and stress, buckling load and frequency responses are selected as design variables in different applications. In this section, exact solutions for the vibration behavior of CNTRC beams

with arbitrary boundary conditions are developed which can be used in the conceptual design.

2.1 Nonlocal elasticity theory

The nonlocal theory of elasticity assumes that the stress at a point is related not only to the strain at that point but also to strains at all other points of the body. In the non-local elasticity, the uniaxial constitutive law is expressed as elasticity. The nonlocal stress tensor at the point is given by Eringen (1983)

$$\sigma^{nl}(x) = \int_{\Omega} K(|x' - x|, \tau) T(x) d\Omega(x) \quad (1a)$$

$$T(x) = C(x) : \varepsilon(x) \quad (1b)$$

where the macroscopic stress tensor at a point “ \mathbf{x} ” is represented by “ T ” and depends on the strain at the same point which is based on the generalized Hook’s law. “ C ” is the fourth-order elasticity tensor which represents the double-dot product. The kernel function “ $K(|x' - x|, \tau)$ ” denotes the nonlocal modulus, “ $(\mathbf{x}' - \mathbf{x})$ ” indicates the distance and “ τ ” is the material constant which depends on the type of material, “ $\varepsilon(x)$ ” is the strain tensor. “ τ ” is the material constant which is defined as “ $\tau = \frac{e_0 \alpha}{l}$ ” and where “ e_0 ” is a constant appropriate to each material, “ α ” is an internal characteristics length and “ l ” is an external characteristics length. In general, a conservative estimate of the nonlocal parameter is “ $e_0 \alpha < 2.0$ nm” for a single-walled CNT (Tounsi *et al.* 2013). Since the solution of the constituent equation in integral form is complex, a simplified differential form equation can be used as given below

$$L\sigma^{nl}(x) = (1 - \mu \nabla^2)\sigma^{nl} = T \quad (2)$$

$$\mu = \tau^2 l^2$$

where “ L^2 ” is the Laplacian operator and “ $L = (1 - \mu \nabla^2)$ ” indicates nonlocal differential. The nonlocal parameter which is represented by “ μ ” varies in accordance with different materials. The nonlocal behavior can be neglected in the thickness direction for a beam type structure.

The nonlocal constitutive relation takes the following form based on classical beam theory

$$\sigma_{xx} - \mu \frac{\partial^2 \sigma_{xx}}{\partial x^2} = E(z)\varepsilon_{xx} \quad (3)$$

2.2 Governing equations

The displacement field of classical beam theory can be expressed as

$$\begin{aligned} u(x, z, t) &= u_0(x, t) - z \frac{\partial w_0}{\partial x} \\ v(x, z, t) &= 0 \\ w(x, z, t) &= w_0 \end{aligned} \quad (4)$$

The strain field of the beam is obtained as

$$V_{NP} = \frac{W_{NP}}{W_{NP} + (\rho_{NP}/\rho_M) - (\rho_{NP}/\rho_M)W_{NP}}, \quad (10)$$

where W_{NP} is the weight fraction of nanoparticles. The potential and kinetic energies of the structure are:

$$\varepsilon_{xx} = \varepsilon_{xx}^0 - z\kappa_x \quad (5)$$

where

$$\varepsilon_{xx}^0 = \frac{\partial u}{\partial x}; \quad \kappa_x = \frac{\partial^2 w_0}{\partial x^2} \quad (6)$$

To obtain the equations of motion, the expression for the Hamilton's principle of a CNTRC beam can be written as

$$\int_{T(\text{time})} (\delta U + \delta V - \delta K) dT = 0 \quad (7)$$

where “ δU ” is the virtual variation of the total strain energy, “ δV ” is the virtual work and “ δK ” is the virtual kinetic energy, which are obtained as

$$\begin{aligned} \delta U &= \int_0^L \int_A \sigma_{xx} \delta \varepsilon_{xx} dA dx \\ &= \int_0^L (N_x \delta \varepsilon_{xx}^0 - M_x \delta \kappa_x) dx \end{aligned} \quad (8)$$

$$\delta V = - \int_0^L \left(q \delta w_0 + N_{x0} \frac{\partial w_0}{\partial x} \delta \frac{\partial w_0}{\partial x} \right) dx \quad (9)$$

$$\begin{aligned} \delta K &= \int_0^L \int_A \rho(z) [u \delta u + w \delta w] dA dx \\ &= \int_0^L \left[I_0 (u_0 \delta u_0 + w_0) + I_1 \left(u_0 \frac{\partial \delta w_0}{\partial x} + \frac{\partial w_0}{\partial x} \delta u_0 \right) \right. \\ &\quad \left. + I_2 \frac{\partial w_0}{\partial x} \delta \frac{\partial w_0}{\partial x} \right] dx \end{aligned} \quad (10)$$

The mass moment of inertia can be represented as

$$\begin{Bmatrix} I_0 \\ I_1 \\ I_2 \end{Bmatrix} = \int_A \rho(z) \begin{Bmatrix} 1 \\ z \\ z^2 \end{Bmatrix} dA \quad (11)$$

The “ N_x ” is the force resultant and the “ M_x ” is the moment resultant that they are defined as below

$$N_x = \int_A \sigma_{xx} dA \quad (12a)$$

$$M_x = \int_A z \sigma_{xx} dA \quad (12b)$$

Substituting Eqs. (8)-(10) into Eq. (7), integrating by parts and setting the coefficient “ δu_0 ” and “ δw_0 ” to zero lead to the following Euler-Lagrange equations:

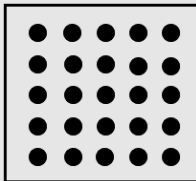
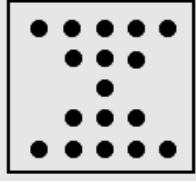
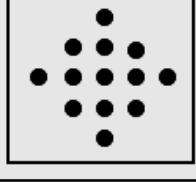
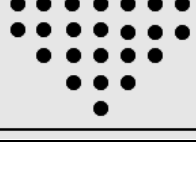
$$\frac{\partial N}{\partial x} + f = 0 \quad (13a)$$

$$\frac{\partial^2 M_x}{\partial x^2} - q = 0 \quad (13b)$$

By using Eqs. (3), (5) and (12), the force-strain and the moment-strain relations of the nonlocal classical beam theory can be obtained as

$$N - \mu \frac{\partial^2 N}{\partial x^2} = A_{xx} \frac{\partial u}{\partial x} - B_{xx} \frac{\partial^2 w}{\partial x^2} \quad (14a)$$

Table 1 Volume fraction and distribution forms of CNT

Distribution form	Volume fraction
	UD: $V_{CNT}(z) = V_{CNT}^*$
	XD: $V_{CNT}(z) = \frac{4 z }{h} V_{CNT}^*$
	OD: $V_{CNT}(z) = 2 \left(1 - \frac{2 z }{h} \right) V_{CNT}^*$
	VD: $V_{CNT}(z) = \left(1 + \frac{2 z }{h} \right) V_{CNT}^*$

$$M_x - \mu \frac{\partial^2 M_x}{\partial x^2} = B_{xx} \frac{\partial u}{\partial x} - D_{xx} \frac{\partial^2 w}{\partial x^2} \quad (14b)$$

where the following cross-sectional rigidities are defined as

$$(A_{xx}, B_{xx}, D_{xx}) = \int_A E(z) (1, z, z^2) dA \quad (15)$$

2.3 CNTRC micro-mechanic model

The effective material properties of the two phase composites, the mixture of CNTs and an isotropic polymer, can be determined from the rule of the mixture as

$$E_{11} = \eta_1 V_{CNT} E_{11}^{CNT} + V_m E_m \quad (16)$$

$$\frac{\eta_2}{E_{22}} = \frac{V_{CNT}}{E_{22}^{CNT}} + \frac{V_m}{E_m} \quad (17)$$

$$\frac{\eta_3}{G_{12}} = \frac{V_{CNT}}{G_{12}^{CNT}} + \frac{V_m}{G_m} \quad (18)$$

$$v_{12} = V_{CNT} v_{12}^{CNT} + V_m v_m \quad (19)$$

$$\rho = V_{CNT} \rho_{CNT} + V_m \rho_m \quad (20)$$

where “ E_{11}^{CNT} ”, “ E_{22}^{CNT} ” and “ G_{12}^{CNT} ” are Elasticity modulus and shear modulus of CNT, “ E_m ” and “ G_m ” are elasticity and shear modulus of the isotropic matrix, “ η_1 ”, “ η_2 ” and “ η_3 ” are CNT/matrix efficiency parameters, “ v_{12}^{CNT} ” and “ v_m ” are Poisson's ratios of CNT and matrix, “ ρ_{CNT} ” and “ ρ_m ” are mass densities of CNT and matrix, respectively. Also,

“ V_{CNT} ” and “ V_m ” are volume fractions of CNT and matrix, respectively and related by

$$V_{CNT} + V_m = 1 \tag{21}$$

where “ V_{CNT}^* ” is the volume fraction of CNT that is determined by

$$V_{CNT}^* = \frac{w_{CNT}}{w_{CNT} + \frac{\rho_{CNT}}{\rho_m} - \frac{\rho_{CNT}}{\rho_m} w_{CNT}} \tag{22}$$

where “ w_{CNT} ” is the mass fraction of CNT.

The nonlocal constitutive relations for the k th layer of the CNTRC beam take the form

$$\sigma_{xx}^{(k)} - \mu \frac{\partial^2 \sigma_{xx}^{(k)}}{\partial x^2} = \bar{Q}_{11}^{(k)}(z) \epsilon_{xx}^{(k)} \tag{23}$$

where

$$\begin{aligned} \bar{Q}_{11}^{(k)} &= Q_{11}^{(k)}(z) \cos^4 \theta \\ &+ 2 \left(Q_{12}^{(k)}(z) + 2Q_{66}^{(k)}(z) \right) \sin^2 \theta \cos^2 \theta \\ &+ Q_{22}^{(k)}(z) \sin^4 \theta \end{aligned} \tag{24}$$

$$\begin{aligned} Q_{11}^{(k)}(z) &= \frac{E_{11}^{(k)}(z)}{1 - \nu_{12}^{(k)}(z)\nu_{21}^{(k)}(z)}; \\ Q_{21}^{(k)}(z) &= \frac{\nu_{12}^{(k)}(z)E_{22}^{(k)}(z)}{1 - \nu_{12}^{(k)}(z)\nu_{21}^{(k)}(z)} \\ Q_{22}^{(k)}(z) &= \frac{E_{22}^{(k)}(z)}{1 - \nu_{12}^{(k)}(z)\nu_{21}^{(k)}(z)}; \\ Q_{66}^{(k)}(z) &= G_{12}^{(k)}(z) \end{aligned} \tag{25}$$

These parameters depend on the CNT volume fraction of the k th layer.

2.4 Analytical solution

The nonlocal axial normal force and bending moment can be written as

$$N = A_{xx} \frac{\partial u}{\partial x} - B_{xx} \frac{\partial^2 w}{\partial x^2} - \mu \frac{\partial^2 f}{\partial x^2} \tag{26a}$$

$$M_x = B_{xx} \frac{\partial u}{\partial x} - D_{xx} \frac{\partial^2 w}{\partial x^2} + \mu \left(\frac{\partial^2 w}{\partial x^2} \bar{N} - q \right) \tag{26b}$$

The nonlocal governing equations in terms of the displacements can be obtained by substituting for N , M from Eqs. (26a, b), respectively, into Eqs. (14a, b) as follows:

$$A_{xx} \frac{\partial^2 u}{\partial x^2} - B_{xx} \frac{\partial^3 w}{\partial x^3} = \mu \frac{\partial^2 f}{\partial x^2} - f \tag{27a}$$

$$B_{xx} \frac{\partial^3 u}{\partial x^3} - D_{xx} \frac{\partial^4 w}{\partial x^4} - \bar{N} \frac{\partial^2 w}{\partial x^2} + \mu \bar{N} \frac{\partial^2 w}{\partial x^2} = \mu \frac{\partial^2 q}{\partial x^2} - q \tag{27b}$$

The exact solution of natural vibration of CNTRC beam was considered and the boundary conditions are

$$SS \rightarrow w_{x=0} = 0 \quad M_{x=L} = 0 \tag{28}$$

$$CC \rightarrow w_{x=0} = 0$$

$$CF \rightarrow w_{x=0} = 0 \quad M_{x=L} = 0$$

$$CS \rightarrow w_{x=L} = 0 \quad M_{x=L} = 0$$

The following expansions of the generalized displacements satisfy the boundary conditions

$$W(x, t) = \sum_{n=1}^{\infty} W_n \sin \frac{n\pi x}{L} e^{i\omega_n t} \tag{29}$$

The general solution is

$$W(x) = c_1 \sin \lambda x + c_2 \cos \lambda x + c_3 \sinh \lambda x + c_4 \cosh \lambda x \tag{30}$$

where “ c_1 , c_2 , c_3 and c_4 ” are integration constants, which must be evaluated using the boundary conditions.

The natural frequencies are given by

$$\begin{aligned} \omega^2 &= \frac{1}{\lambda M_x} \left(\frac{\eta\pi}{L} \right)^4 E(z) I_{xx} \\ \lambda &= 1 + \mu \left(\frac{n\pi}{L} \right)^2 \end{aligned} \tag{31}$$

For free vibration, we set “ \bar{N} ” and “ q ” to zero.

2.5 Artificial Neural Networks (ANNs)

Artificial Neural Networks (ANNs) is a technology that is inspired by the biological nervous system and performs the learning feature, which is the most important feature of the human brain. The learning feature refers to the ability of a neural network to make certain inferences as a result of its experiences. ANN learns from examples and makes decisions about the examples they encounter later using their previous knowledge. Thus, it makes generalizations about new examples. ANN has been successfully used in many different areas such as agriculture (Eldem 2020), civil engineering (Madenci and Gülcü 2020), chemical engineering (Tümer *et al.* 2020), and many different problems such as the classification of type 2 diabetes (Frimpong *et al.* 2021), human gender prediction (Hacibeyoglu and Ibrahim 2018), and optimization of traffic signal timing (Karaşahin and Tümer).

Many ANN models have been proposed in the literature. A few of them are feed-forward artificial neural network (FFNN) (Bebis and Georgiopoulos 1994), Kohonen self-organizing network (Kohonen 1982), radial basis function network (Park and Sandberg 1993), recurrent neural network (Dorffner 1996). The most widely used model of these is the FFNN model.

FFNN consists of a series of data processing elements called neurons. Neurons are connected to each other and are located in units called layers. The first of these layers is called the input layer. The input layer maps the input data of the problem to the network by passing it on to the neurons. The last layer in the network structure is called the output layer. The layer between the input layer and the output layer is called the hidden layer. The number of hidden layers in the network and the number of neurons in the layers may differ depending on the difficulty level of the problem. The

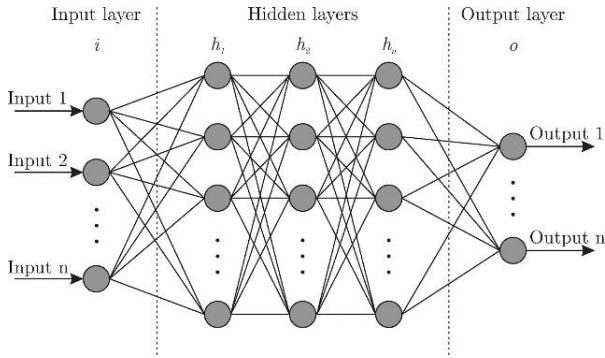


Fig. 1 The general architecture of the FFNN (Bre *et al.* 2018)

Table 2 Dataset properties used in the experiments

CNT Distribution model	Boundary Conditions	Fiber Angles
UD	SS	0°
XD	CC	15°
OD	CF	30°
VD	SC	45°
		60°
		75°
		90°

Table 3 Statistical analysis of data

Attribute	Min	Max	Mean	St. Dev.
CNT Distribution model	1	4	2.5	1.12
Boundary Conditions	1	4	2.5	1.12
Fiber Angles	1	7	4	2.01
Natural frequencies	0.0716	1.5147	0.6458	0.44

connections that allow neurons to connect to each other are called weights.

The general architecture of the FFNN is shown in Fig. 1. The outputs of neurons in a layer are given as inputs to the neurons in the next layer by passing through the weights. Information in the network is transferred in one direction from the input layer to the output layer. The performance of FFNN depends on the updating of the network with optimum values of the weights and biases.

The output of the FFNN with one hidden layer is calculated as follows (Mirjalili 2015, Gülcü 2020): The inputs, the initial values of weights, and biases are given to the network. In the first stage, the weighted sum value of the input is calculated by multiplying the input data and the connection weights between the input layer and hidden layer. Eq. (32) shows the weighted sum calculation.

$$t_i = \sum_{k=1}^m (A_{ki} * G_k) - \beta_i \quad i = 1, 2, \dots, n \quad (32)$$

where “ m , n , k , and i ” represent the total number of neurons in the input layer, the total number of neurons in the hidden layer, a neuron in the input layer, and a neuron in the hidden layer, respectively. “ A_{ki} ” represents the weight between the

k th neuron in the input layer and the “ i th” neuron in the hidden layer. “ G_k ” represents the input data. “ β_i ” represents the bias value of the “ i th” neuron in the hidden layer. The output of each neuron in the hidden layer is calculated by the sigmoid function shown in Eq. (33).

$$T_i = Sigmoid(t_i) = \frac{1}{1 + e^{-t_i}} \quad (33)$$

where i and t_i represent a neuron in the hidden layer and the weighted sum calculation of the i th neuron, respectively. Eqs. (34) and (35) are used to calculate the actual output of the network.

$$c_y = \sum_{k=1}^m (A_{iy} * T_i) - \beta_y \quad y = 1, 2, \dots, p \quad (34)$$

$$C_y = Sigmoid(c_y) = \frac{1}{1 + e^{-c_y}} \quad (35)$$

where n , p , i , and y represent the total number of neurons in the hidden layer, the total number of neurons in the output layer, a neuron in the hidden layer, and a neuron in the output layer, respectively. A_{iy} represents the weight between i th neuron in the hidden layer and y th neuron in the output layer. T_i represents the output of the i th neuron in the hidden layer. β_y represents the bias value of the y th neuron in the output layer. c_y represents the weighted sum calculation of the y th neuron. C_y represents the output of the network. There are six fundamental steps in the application of FFNN: collecting and preparing data, defining network architecture, initializing weights and biases, training network, validating network, and using the network.

Collecting and preparing data is the first step of the application of the FFNN. Therefore, data should be coherent and should not contain missing values. To create a good neural network, 112 different design mixed data-set was obtained based on free vibration analysis of CNTRC layered beam. Where the input data are CNT distribution models, boundary conditions and fiber angles; output data is natural frequencies, respectively. The combination parameters of input data are given in Table 2.

The statistical analysis of the data is shown in Table 3.

In the step of preparing data, the instances of the training dataset, the validation dataset, and the test dataset are randomly determined. In the application of FFNN, a common task is the construction of a suitable model that can learn from data and make predictions on data. The data used to build the final model usually comes from multiple datasets. In particular, three datasets are commonly used in different stages of the creation of the model. They are the training dataset, the validating dataset, and the test dataset. The training dataset is a set of examples used to fit the weights and biases of the model. FFNN is trained using the training dataset to create the fittest model. Successively, the fitted model is used to predict the responses for the observations in a second dataset called the validation dataset (James *et al.* 2013). Finally, the test dataset is used to provide an unbiased evaluation of a final model fit on the training dataset. In general, the training data, the validation data, and the test data are 70%, 15%, and 15% of the data, respectively. In this paper, we also used the same

fragmentation. In addition, all data sets are normalized by using the min-max normalization function given in Eq. (36) to eliminate the effect of attributes that may have different effective rates on the classification (Aljarah *et al.* 2018).

$$x' = \frac{(x - x_{min})}{(x_{max} - x_{min})} \quad (36)$$

where x' represents the normalized form of x .

An FFNN contains one input layer, one or many hidden layers, and one output layer. In this study, we used the architecture of the FFNN with one hidden layer. Each layer contains one or many neurons depends on the problem. The input layer and the output layer have three neurons and one neuron respectively since the dataset used in the experiments contains three input and one output. There is no fixed rule in determining the number of neurons in the hidden layer. The number of neurons in the hidden layer is empirically determined using the validation dataset. In this study, the number of hidden neurons was started at two and was tried up to ten. Therefore, nine different FFNN architectures were used to determine the best model. The values of the weights and biases affect the performance of the FFNN. The values of the weights and biases are randomly initialized at the beginning of the algorithm. There is no fixed rule in initializing the values of the weights and biases. But they are often initialized in the range [-0.5, 0.5] or [-1, 1]. Therefore, they are initialized between -1 and 1 in this study.

In the training network, FFNN is trained using the training dataset which is 70% of the data in this study. The weights and biases of FFNN are optimized according to the inputs-outputs pattern in the training dataset. In this study, the mean square error (MSE) was used as the objective function whose formulation is given in Eq. (37). The aim is to minimize the MSE. In this study, the number of hidden neurons was started at two and was tried up to ten. Therefore, nine different FFNN architectures were used to determine the best model. At the end of the training network phase, nine models with the minimum MSE are created. Successively, the best model among nine models is found using the validation dataset. Finally, the test dataset is used to provide an unbiased evaluation of the best model fit on the training dataset. Thus, the best model is ready to use.

To evaluate the performance of FFNN, the statistical metrics which are frequently used in literature were applied. These statistical metrics are the MSE, the correlation of the coefficient (R), and the mean absolute error (MAE). Their formulations are shown in Eqs. (37)-(39).

$$MSE = \frac{(\sum_{i=1}^N (O_i - P_i)^2)}{N} \quad (37)$$

$$MAE = \frac{(\sum_{i=1}^N |O_i - P_i|)}{N} \quad (38)$$

$$R = \sqrt{1 - \frac{(\sum_{i=1}^N (O_i - P_i)^2)}{(\sum_{i=1}^N (O_i - \bar{P})^2)}} \quad (39)$$

where N , O , P , and \bar{P} are the number of instances in the dataset, the observed value, the predicted value, and the mean of the predicted values, respectively.

Table 4 Material parameters of a CNTRC beam

Parameters	CNT	Matrix
Young's modulus	$E_{11} = 600$ GPa, $E_{22} = 10$ GPa	$E = 2.5$ GPa
Poisson's coefficient	$\nu = 0.3$	$\nu_{12} = 0.19$
mass density	$\rho = 1400$ kg/m ³	$\rho = 1190$ kg/m ³
shear modulus	$G_{12} = 17.2$ GPa	---

Table 5 The CNT efficiency parameters

V_{CNT}^*	η_1	η_2	η_3
0.12	1.2833	1.0556	1.0556
0.17	1.3414	1.7101	1.7101
0.28	1.3238	1.7380	1.7380
shear modulus	$G_{12} = 17.2$ GPa	---	---

3. Numerical results

For the first comparison study, the fundamental frequency parameter of the layered CNTRC beam is evaluated and compared with literature results. The material properties of the CNTRC beam are given in Table 4.

The CNT efficiency parameters " η_i " associated with the given volume fraction " V_{CNT}^* " are shown in Table 5. It should be mentioned that the UD and other FG-CNT distributions nanobeams have the same mass fraction of CNT.

The non-dimensional parameter is defined as

$$\bar{\omega} = \omega \frac{L^2}{h} \sqrt{\frac{\rho_m}{E_m}} \quad (40)$$

Tables 6-9 provide the non-dimensional natural frequencies of four-layered simply-supported (SS) angleply CNTRC beams by the classical and nonlocal theories are compared for different distributions of CNTs for $V_{CNT}^*=0.12$ case. Frequency values started to decrease with the increase of nonlocal parameter. Considering the effect of the layer angle, the frequency value decreased with the increase of the angle. CNT distribution had an effect on frequency. While the lowest frequency value was in the FG-O distribution, the highest frequency value occurred in the FG-X distribution. The results are also presented for different boundary condition cases. Tables 10-13 give the results for the clamped-clamp (CC) boundary condition. Compared to the SS support, the CC support results were lower. Frequency values decreased with the increase of nonlocal parameter and angle. The frequency values of the cantilever beam are given in Tables 14-17. Due to the boundary conditions, the frequency values decreased significantly. The frequency results of beams with CS boundary condition are presented in Tables 18-21. The frequency values for the CS beam were close to the results of the CC beam. In addition, the results were compared with the mixed finite element solution results made by Madenci (2021) and the local results where the u parameter was 0, and the results were the same.

Table 6 The non-dimensional natural frequencies of SS CNTRC beams for UD distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	0.9748	0.9281	0.7886	0.5715	0.3417	0.2139	0.2003
2	0.8899	0.8473	0.7199	0.5217	0.3119	0.1952	0.1828
4	0.8133	0.7743	0.6579	0.4768	0.2850	0.1784	0.1671
(Madenci 2021)	0.9748	0.9281	0.7886	0.5715	0.3417	0.2139	0.2003

Table 7 The non-dimensional natural frequencies of SS CNTRC beams for FG-X distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	0.9851	0.9382	0.7976	0.5782	0.3452	0.2149	0.2009
2	0.8993	0.8565	0.7282	0.5278	0.3151	0.1962	0.1834
4	0.8219	0.7827	0.6654	0.4824	0.2880	0.1793	0.1676
(Madenci 2021)	0.9851	0.9382	0.7976	0.5782	0.3452	0.2149	0.2009

Table 8 The non-dimensional natural frequencies of SS CNTRC beams for FG-V distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	0.9733	0.9266	0.787	0.5703	0.3413	0.2142	0.2006
2	0.8886	0.8459	0.7185	0.5206	0.3116	0.1955	0.1831
4	0.8120	0.7730	0.6566	0.4757	0.2847	0.1786	0.1673
(Madenci 2021)	0.9733	0.9266	0.787	0.5703	0.3413	0.2142	0.2006

Table 9 The non-dimensional natural frequencies of SS CNTRC beams for FG-O distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	0.9639	0.9174	0.7785	0.5638	0.3381	0.2135	0.2003
2	0.8800	0.8375	0.7107	0.5147	0.3086	0.1949	0.1828
4	0.8041	0.7654	0.6495	0.4704	0.2821	0.1781	0.1671
(Madenci 2021)	0.9639	0.9174	0.7785	0.5638	0.3381	0.2135	0.2003

Table 10 The non-dimensional natural frequencies of CC CNTRC beams for UD distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	1.5052	1.4710	1.3507	1.0985	0.7265	0.4729	0.4441
2	1.3742	1.3430	1.2331	1.0029	0.6632	0.4317	0.4054
4	1.2558	1.2273	1.1269	0.9165	0.6061	0.3945	0.3705
(Madenci 2021)	1.5052	1.4710	1.3507	1.0985	0.7265	0.4729	0.4441

Table 11 The non-dimensional natural frequencies of CC CNTRC beams for FG-X distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	1.5147	1.4807	1.361	1.1086	0.7332	0.4749	0.4454
2	1.3829	1.3518	1.2425	1.0121	0.6694	0.4335	0.4066
4	1.2637	1.2354	1.1355	0.9249	0.6117	0.3962	0.3716
(Madenci 2021)	1.5147	1.4807	1.361	1.1086	0.7332	0.4749	0.4454

Table 12 The non-dimensional natural frequencies of CC CNTRC beams for FG-V distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	1.5024	1.4677	1.346	1.0928	0.7235	0.4733	0.4448
2	1.3716	1.3400	1.2288	0.9977	0.6605	0.4321	0.4061
4	1.2535	1.2245	1.1230	0.9117	0.6036	0.3948	0.3711
(Madenci 2021)	1.5024	1.4677	1.346	1.0928	0.7235	0.4733	0.4448

Table 13 The non-dimensional natural frequencies of CC CNTRC beams for FG-O distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	1.4997	1.4647	1.3425	1.0888	0.7199	0.4721	0.4442
2	1.3692	1.3372	1.2257	0.9940	0.6572	0.4310	0.4055
4	1.2512	1.2220	1.1201	0.9084	0.6006	0.3938	0.3706
(Madenci 2021)	1.4997	1.4647	1.3425	1.0888	0.7199	0.4721	0.4442

Table 14 The non-dimensional natural frequencies of CF CNTRC beams for UD distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	0.3761	0.3553	0.2956	0.209	0.1229	0.0765	0.0716
2	0.3433	0.3243	0.2698	0.1908	0.1122	0.0698	0.0653
4	0.3138	0.2964	0.2466	0.1743	0.1025	0.0638	0.0597
(Madenci 2021)	0.3761	0.3553	0.2956	0.209	0.1229	0.0765	0.0716

Table 15 The non-dimensional natural frequencies of CF CNTRC beams for FG-X distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	0.3806	0.3596	0.2992	0.2115	0.1241	0.0768	0.0718
2	0.3474	0.3283	0.2731	0.1930	0.1133	0.0701	0.0655
4	0.3175	0.3000	0.2496	0.1764	0.1035	0.0640	0.0599
(Madenci 2021)	0.3806	0.3596	0.2992	0.2115	0.1241	0.0768	0.0718

Table 16 The non-dimensional natural frequencies of CF CNTRC beams for FG-V distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	0.3727	0.3520	0.2928	0.2071	0.1222	0.0766	0.0717
2	0.3402	0.3213	0.2673	0.1890	0.1115	0.0699	0.0654
4	0.3109	0.2936	0.2442	0.1727	0.1019	0.0639	0.0598
(Madenci 2021)	0.3727	0.3520	0.2928	0.2071	0.1222	0.0766	0.0717

Table 17 The non-dimensional natural frequencies of CF CNTRC beams for FG-O distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	0.3710	0.3504	0.2913	0.2060	0.1215	0.0763	0.0716
2	0.3387	0.3199	0.2659	0.1880	0.1109	0.0696	0.0653
4	0.3095	0.2923	0.2430	0.1718	0.1013	0.0636	0.0597
(Madenci 2021)	0.3710	0.3504	0.2913	0.2060	0.1215	0.0763	0.0716

Table 18 The non-dimensional natural frequencies of CS CNTRC beams for UD distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	1.2426	1.2021	1.0690	0.8250	0.5185	0.3305	0.3098
2	1.1344	1.0975	0.9759	0.7532	0.4733	0.3017	0.2828
4	1.0367	1.0029	0.8919	0.6883	0.4326	0.2757	0.2584
(Madenci 2021)	1.2426	1.2021	1.0690	0.8250	0.5185	0.3305	0.3098

Table 19 The non-dimensional natural frequencies of CS CNTRC beams for FG-X distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	1.2525	1.2120	1.0789	0.8336	0.5235	0.3319	0.3107
2	1.1435	1.1065	0.9850	0.7610	0.4779	0.3030	0.2836
4	1.0450	1.0112	0.9001	0.6955	0.4367	0.2769	0.2592
(Madenci 2021)	1.2525	1.2120	1.0789	0.8336	0.5235	0.3319	0.3107

Table 20 The non-dimensional natural frequencies of CS CNTRC beams for FG-V distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	1.2390	1.1981	1.0642	0.8203	0.5164	0.3308	0.3103
2	1.1312	1.0938	0.9716	0.7489	0.4714	0.3020	0.2833
4	1.0337	0.9996	0.8879	0.6844	0.4308	0.2760	0.2589
(Madenci 2021)	1.2390	1.1981	1.0642	0.8203	0.5164	0.3308	0.3103

Table 21 The non-dimensional natural frequencies of CS CNTRC beams for FG-O distributions

μ	0°	15°	30°	45°	60°	75°	90°
0	1.2344	1.1934	1.0594	0.8158	0.5133	0.3299	0.3098
2	1.1270	1.0895	0.9672	0.7448	0.4686	0.3011	0.2828
4	1.0299	0.9957	0.8839	0.6806	0.4282	0.2752	0.2584
(Madenci 2021)	1.2344	1.1934	1.0594	0.8158	0.5133	0.3299	0.3098

In the second section, we present the experimental results of FFNN. The specifications of the hardware and software used in the experiments are as follows: Windows 10 operating system, Intel(R) Core(TM) i5-3330 3.00 GHz, 4 GB of memory, visual studio 2017, C# programming language. In this study, the FFNN architecture has three layers (one input, one hidden, and one output). The numbers of neurons in the input and output layers are three and one, respectively. As mentioned above, the number of neurons in the hidden layer was started at two and was tried up to ten. Therefore, nine different FFNN architectures were used to determine the best model. The best model was found using the validation dataset and had nine neurons in the hidden layer. The best model has the architecture with 3-9-1. The momentum factor, the learning rate, and the maximum epoch number are 0.8, 0.3, and 1000, respectively. Eq. (41) shows the formulation of the generated best model. Table 22 shows the values of weights in the hidden layer of the generated best model. Table 23 shows the values of the weights and the biases in the output layer of the generated

Table 22 The weights of the hidden layer

IHW		
0.1377	-13.8239	-0.4561
0.3644	-3.7677	2.5349
-0.0021	-3.2193	-2.9911
-0.3041	-1.8542	-2.4603
-1.0881	-2.2786	0.4406
0.6180	-2.9044	-0.1490
0.4843	0.4823	-2.4643
0.0049	0.1872	5.8122
-0.2262	5.2330	1.4390

Table 23 The biases in the hidden layer and the weights and the biases in the output layer

b_1	HOW	b_2
6.7103	9.0526	-1.2402
0.1532	-3.0927	
-0.1934	-3.4070	
-0.4038	-2.0494	
0.8941	-2.3713	
0.0879	-2.9252	
-0.4825	0.1584	
-2.9276	-4.1444	
-4.4383	5.6203	

Table 24 Performance of the generated FFNN model

Criteria	Training	Validation	Test
Instance	78	17	17
R	0.996	0.993	0.996
MAE	2.20E-02	2.72E-02	2.53E-02
MSE	7.18E-04	9.72E-04	9.08E-04

best model.

$$Y = sigmoid(sigmoid(X \times IHW + b_1) \times HOW + b_2) \tag{41}$$

where $sigmoid(\cdot)$ is the activation function, X represents the inputs. IHW and HOW represent the values of the weights in the hidden layer and output layer, respectively. b_1 and b_2 represent the values of the biases in the hidden layer and output layer, respectively.

Table 24 shows the performance results of the generated model for three sub-datasets. There are 112 instances in the dataset. The dataset is divided into 70% for training, 15% for validation, and 15% for test using stratified sampling in order to protect the class distribution as much as possible. Namely, 78 instances for training, 17 instances for validation, and 17 instances for test. Moreover, to eliminate the effect of the attributes in the dataset that have different scales, it is normalized using Eq. (36). The values of the correlation of the coefficient (R) for the training, validation, and test are 0.996, 0.9993, and 0.996, respectively. The

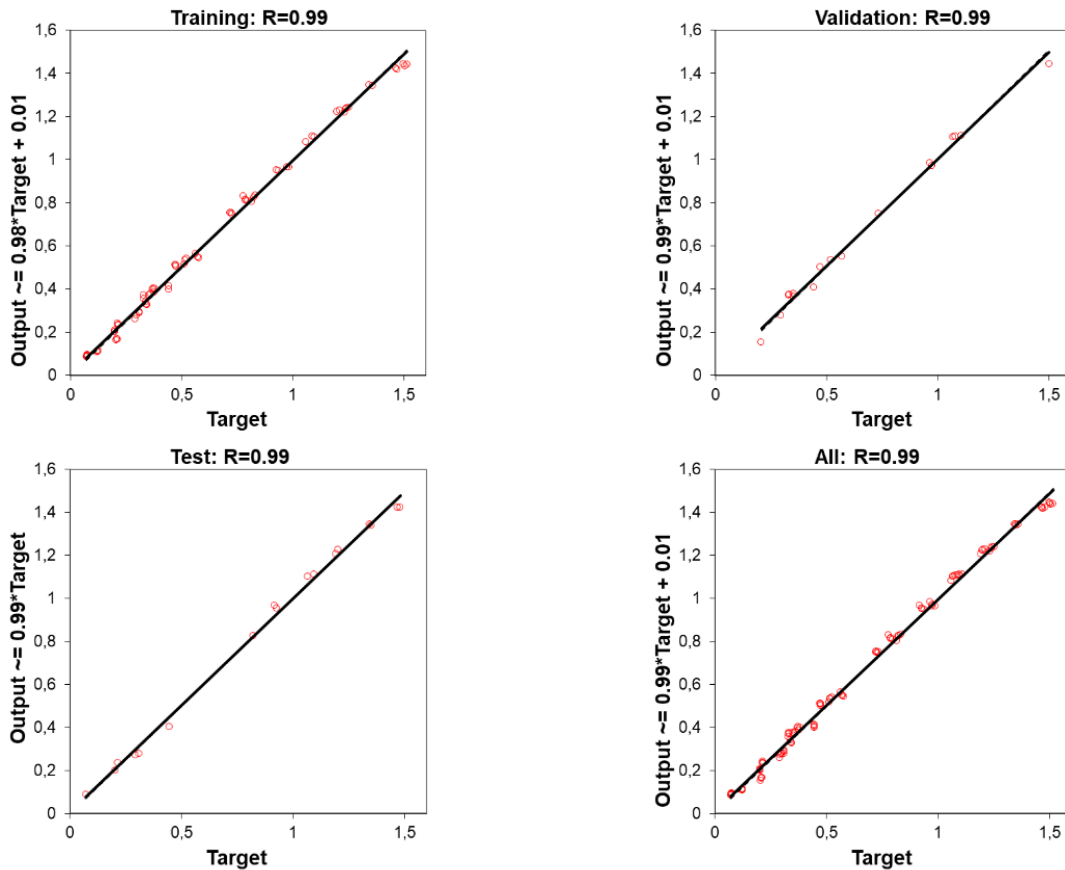


Fig. 2 The regression results of the training, validation, test, and all data

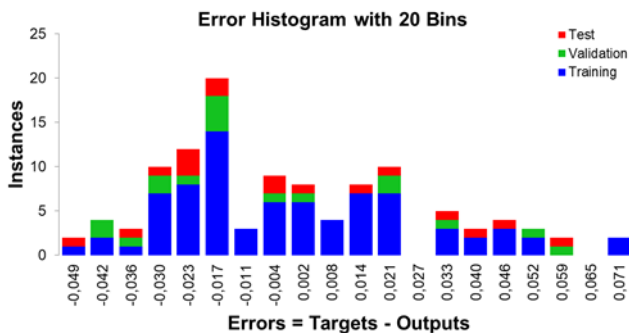


Fig. 3 The error histogram

larger the coefficient, the stronger the relationship, so that a correlation that is close to one indicates a very strong relationship, while coefficients that are near zero indicate very weak relationships (Bachman 2004). Because the values of the R for the training, validation, and test are very close to one, there is a very strong relationship. The values of the MAE for the training, validation and test are 2.20E-02, 2.72E-02, and 2.53E-02, respectively. The values of the MSE for the training, validation, and test are 7.18E-04, 9.72E-04, and 9.08E-04, respectively. Lower MSE and lower MAE indicate more accurate predictions. According to these results, the generated FFNN model has a good performance and accurate prediction.

The regression plots of the generated FFNN model are shown in Fig. 2. The prediction performance of the proposed

FFNN model is very good according to the plots of the training, validation, test, and all data.

The error histogram of the generated FFNN model is shown in Fig. 3. It provides additional verification of the performance of the artificial neural network. It indicates outliers. The blue, green, and red areas represent training, validation, and test data respectively. Most data fall on zero error which provides an idea to check the outliers to determine if the data is bad, or if those data points are different than the rest of the data set. If the outliers are valid data points but are unlike the rest of the data, then the network is extrapolating for these points (Yadav et al. 2015). The density in the error histogram is around zero as shown in Fig. 3. Therefore, the prediction success of the generated FFNN model is very high.

4. Conclusions

In this study, an ANN-based model for the free vibration response of layered CNTRC beams was also presented. Dynamic characteristics of beams made up of FG-CNTRC nanobeams have been illustrated using the nonlocal classical beam theory. Nanobeams were considered to have a combination of FG-CNTRC layers with single-walled CNT. The effective material properties were obtained from the rule of the mixture by taking into account four types of volume distribution models using micromechanics theory.

The material properties of FG- CNTRCs are assumed to be graded through the thickness direction according to several distributions namely UD, FG-X, FG-V and FG-O. The main feature of the ANN is the learning ability from data. The main aim of this study is to develop a model which may predict the natural frequency. The proposed ANN model is composed of one input layer, one hidden layer, and one output layer. The architecture of the model is 3-9-1. In the experiments, we used 112 different data to predict the natural frequency using ANN. The input data in ANN are the CNT distribution models, boundary conditions, and fiber angles. The output data in ANN is the natural frequency. The train data are 70% of the data, the validation data are 15% of the data, and the test data are 15% of the data. The proposed ANN model is evaluated by the correlation of the coefficient (R), the mean absolute error (MAE), and the mean square error (MSE) statistical methods. The values of the R for the training, validation, and test are 0.996, 0.9993, and 0.996, respectively. Additionally, the values of the MAE and MSE are close to zero. In conclusion, the prediction ability of the ANN model is very good and the natural frequency can be predicted in ANN without attempting any experiments.

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